

# $A_{FB}$ in semileptonic decays at Belle II

*Anomalies and Precision in the Belle II Era*

Session on  $b \rightarrow c l \nu$  transitions

Manca Mrvar  
*for the Belle II collaboration*

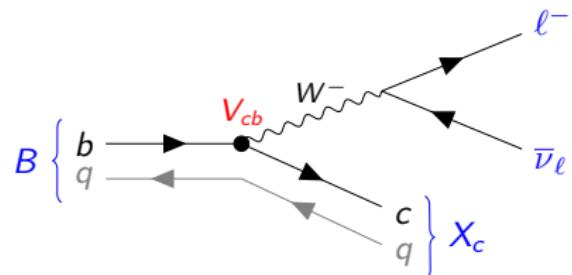
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7th September 2021

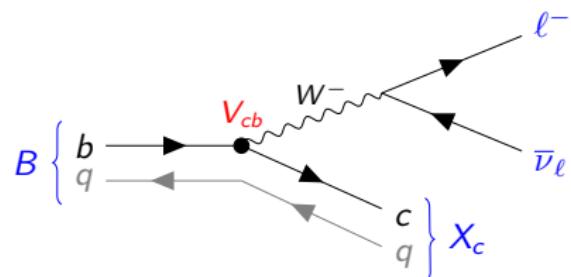
$b \rightarrow c$  transitions: Beyond the  $|V_{cb}|$ 

Probing the  $b \rightarrow c$  transitions with different semileptonic decays:  $B \rightarrow X_c \ell \nu_\ell$ ,  
 $B \rightarrow D^* \ell \nu_\ell$ ,  $B \rightarrow D^0 \ell \nu_\ell$ , etc.



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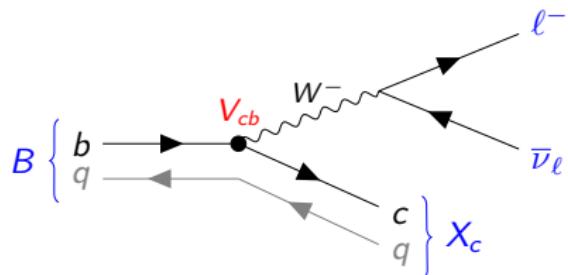


- A well-known anomaly:  $|V_{cb}|$  tension between exclusive and inclusive measurements
- Existence of new physics with no effect on  $|V_{cb}|$ , but on other observables?

# $b \rightarrow c$ transitions: Beyond the $|V_{cb}|$



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- A well-known anomaly:  $|V_{cb}|$  tension between exclusive and inclusive measurements
- Existence of new physics with no effect on  $|V_{cb}|$ , but on other observables?

$$\bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$$

- Large potential: Many observables
- Flavour tagging by slow pion
- Can probe the physics Beyond the Standard Model (BSM)

$$\begin{aligned} \bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell \\ \quad \hookrightarrow D^0 \pi_s^+ \\ \quad \quad \quad \hookrightarrow K^- \pi^+ \end{aligned}$$

# Full angular distribution in $\bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$



[doi:10.1103/PhysRevD.90.074013]

Angular dependency of the semileptonic width:

$$\frac{d^4\Gamma^{(\ell)}}{dq^2 d\cos\theta_\ell d\cos\theta_D d\chi} = \frac{3}{8\pi} \sum_i J_i^{(\ell)}(q^2) f_i(\cos\theta_\ell, \cos\theta_D, \chi) \quad (1)$$

↓ 12 terms

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Two-lepton invariant mass,  $q$ :

$$q^2 = (p_\ell + p_\nu)^2 = \cancel{p_\ell^2}^0 + \cancel{p_\nu^2}^0 + 2(E_\ell E_\nu - \vec{p}_\ell \cdot \vec{p}_\nu) \quad (2)$$

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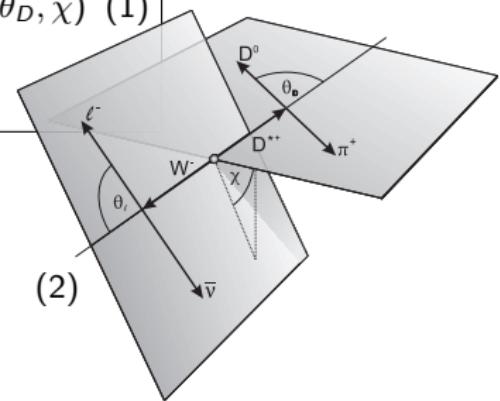


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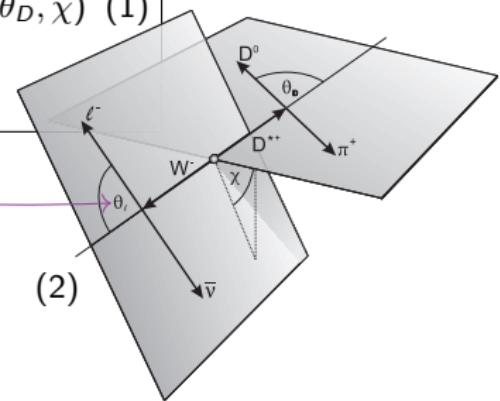
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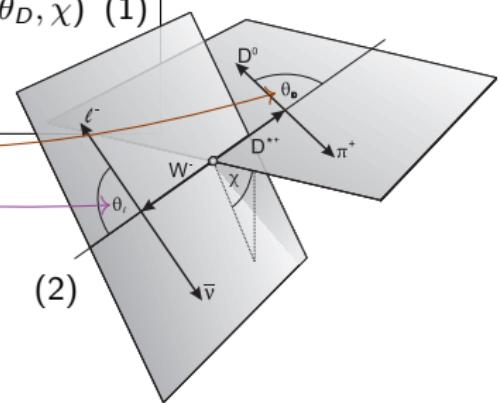
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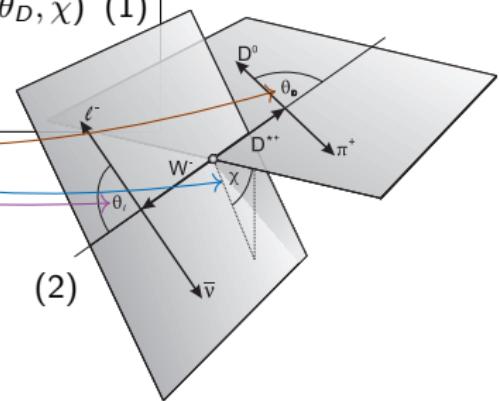
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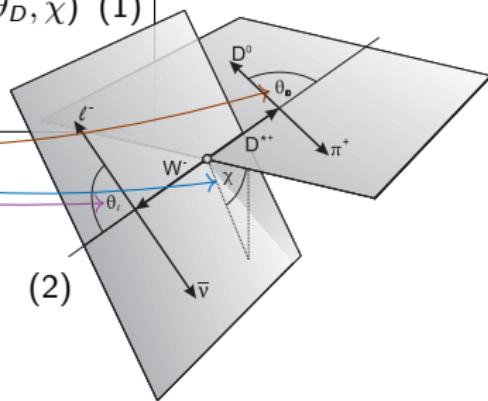


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After integration over one variable,  $q^2$ :

$$\frac{d^3\Gamma^{(\ell)}}{d\cos\theta_\ell d\cos\theta_D d\chi} = \frac{3}{8\pi} \sum_i \langle J_i^{(\ell)} \rangle f_i(\cos\theta_\ell, \cos\theta_D, \chi) \quad (3)$$

→ 12 angular observables  $J_i^{(\ell)}$

→ Their respective (linearly independent) angular coefficient functions  $f_i$

## Angular distributions - cont'd



The number of observables can be reduced by:

- Binned CP-averaging ( $\Gamma \rightarrow \hat{\Gamma}^{(\ell)}$ )
- Integration over angles

# Angular distributions - cont'd



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- Binned CP-averaging ( $\Gamma \rightarrow \hat{\Gamma}^{(\ell)}$ )
- Integration over angles

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d \hat{\Gamma}^{(\ell)}}{d \cos \theta_\ell} = \frac{1}{2} + \left\langle A_{FB}^{(\ell)} \right\rangle \cos \theta_\ell + \frac{1}{4} \left( 1 - 3 \left\langle \tilde{F}_L^{(\ell)} \right\rangle \right) \frac{3 \cos^2 \theta_\ell - 1}{2} \quad (4)$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d \hat{\Gamma}^{(\ell)}}{d \cos \theta_D} = \frac{3}{4} \left( 1 - \left\langle F_L^{(\ell)} \right\rangle \right) \sin^2 \theta_D + \frac{3}{2} \left\langle F_L^{(\ell)} \right\rangle \cos^2 \theta_D \quad (5)$$

$$\frac{1}{\hat{\Gamma}^{(\ell)}} \frac{d \hat{\Gamma}^{(\ell)}}{d \chi} = \frac{1}{2\pi} + \frac{2}{3\pi} \left\langle S_3^{(\ell)} \right\rangle \cos 2\chi + \frac{2}{3\pi} \left\langle S_9^{(\ell)} \right\rangle \sin^2 \chi \quad (6)$$

CP-averaged value vanishes

# Angular distributions - cont'd



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- Binned CP-averaging ( $\Gamma \rightarrow \widehat{\Gamma}^{(\ell)}$ )
- Integration over angles

$$\frac{1}{\widehat{\Gamma}^{(\ell)}} \frac{d \widehat{\Gamma}^{(\ell)}}{d \cos \theta_\ell} = \frac{1}{2} + \left\langle A_{FB}^{(\ell)} \right\rangle \cos \theta_\ell + \frac{1}{4} \left( 1 - 3 \left\langle \tilde{F}_L^{(\ell)} \right\rangle \right) \frac{3 \cos^2 \theta_\ell - 1}{2} \quad (4)$$

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CP-averaged value vanishes ↪

4 independent angular observables that have sensitivity to BSM physics:

$\left\langle A_{FB}^{(\ell)} \right\rangle$  : lepton forward-backward asymmetry

$\left\langle F_L^{(\ell)} \right\rangle$  :  $D^*$  longitudinal polarization factor

$\left\langle \tilde{F}_L^{(\ell)} \right\rangle$  and  $\left\langle S_3^{(\ell)} \right\rangle$  : two further angular observables

# Generic $b \rightarrow c$ model



Effective Field Theory:  $\mathcal{L}(b \rightarrow c l \bar{\nu}_\ell) = \frac{4G_F}{\sqrt{2}} V_{cb} \sum_i \sum_{\ell'} C_i^{\ell\ell'} \mathcal{O}_i^{\ell\ell'} + \text{h.c.}$  (7)

(dim. 6)

Wilson coefficient (unknown) ↗ operator (known) ↘

5 terms ↙ ↘ sum over neutrino flavour

# Generic $b \rightarrow c$ model



Effective Field Theory:  
(dim. 6)

$$\mathcal{L}(b \rightarrow c \ell \bar{\nu}_\ell) = \frac{4G_F}{\sqrt{2}} V_{cb} \sum_i \sum_{\ell'} C_i^{\ell \ell'} \mathcal{O}_i^{\ell \ell'} + \text{h.c.} \quad (7)$$

Wilson coefficients

5 terms

$J_i$  observables

Wilson coefficient (unknown)

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sum over neutrino flavour

Observable	$ C_A ^2$	$ C_V ^2$	$ C_P ^2$	$ C_T ^2$	$\text{Re}(C_A C_V^*)$	$\text{Re}(C_A C_P^*)$	$\text{Re}(C_A C_T^*)$	$\text{Re}(C_V C_P^*)$	$\text{Re}(C_V C_T^*)$	$\text{Re}(C_P C_T^*)$
$J_{1c} = V_1^0$	✓	—	✓	✓	—	(m)	(m)	—	—	—
$J_{1s} = V_1^T$	✓	✓	—	✓	—	—	(m)	—	(m)	—
$J_{2c} = V_2^0$	✓	—	—	✓	—	—	—	—	—	—
$J_{2s} = V_2^T$	✓	✓	—	✓	—	—	—	—	—	—
$J_3 = V_4^T$	✓	✓	—	✓	—	—	—	—	—	—
$J_4 = V_1^{0T}$	✓	—	—	✓	—	—	—	—	—	—
$J_5 = V_2^{0T}$	(m <sup>2</sup> )	—	—	(m <sup>2</sup> )	✓	(m)	(m)	—	(m)	✓
$J_{6c} = V_3^0$	(m <sup>2</sup> )	—	—	—	—	(m)	(m)	—	—	✓
$J_{6s} = V_3^T$	—	—	—	(m <sup>2</sup> )	✓	—	(m)	—	(m)	—
$d\Gamma/dq^2$	✓	✓	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(A_{FB})$	(m <sup>2</sup> )	—	—	(m <sup>2</sup> )	✓	(m)	(m)	—	(m)	✓
$\text{num}(F_L)$	✓	—	✓	✓	—	(m)	(m)	—	—	—
$\text{num}(F_L \cdot 1/3)$	✓	✓	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(\tilde{F}_L)$	✓	(m <sup>2</sup> )	✓	✓	—	(m)	(m)	—	(m)	—
$\text{num}(\tilde{F}_L \cdot 1/3)$	✓	✓	—	✓	—	—	—	—	—	—
$\text{num}(S_3)$	✓	✓	—	✓	—	—	—	—	—	—
Observable	—	—	—	—	$\text{Im}(C_A C_V^*)$	$\text{Im}(C_A C_P^*)$	$\text{Im}(C_A C_T^*)$	$\text{Im}(C_V C_P^*)$	$\text{Im}(C_V C_T^*)$	$\text{Im}(C_P C_T^*)$
$J_7 = V_3^{0T}$					(m <sup>2</sup> )	—	(m)	(m)	—	✓
$J_8 = V_4^{0T}$					✓	—	—	—	—	—
$J_9 = V_5^T$					✓	—	—	—	—	—

TABLE I. The dependence of angular observables on combinations of Wilson coefficients. An entry of ✓ denotes the presence of this combination. A dash means that the observable is not related to this combination.

# Generic $b \rightarrow c$ model



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$$\text{Wilson coefficient (unknown)} \leftarrow \mathcal{L}(b \rightarrow c \ell \bar{\nu}_\ell) = \frac{4G_F}{\sqrt{2}} V_{cb} \sum_i \sum_{\ell'} C_i^{\ell \ell'} \mathcal{O}_i^{\ell \ell'} + \text{h.c.} \quad (7)$$

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$d\Gamma/dq^2$	✓	✓	✓	✓	—	(m)	(m)	—	(m)	—
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TABLE I. The dependence of angular observables on combinations of Wilson coefficients. An entry of ✓ denotes the presence of this combination. A dash means that the combination is not present.

# Belle analysis of $\bar{B} \rightarrow D^* \ell \bar{\nu}_\ell$



Why Belle (II)?  $B \rightarrow D^*$  angular analysis needs high-precision measurements!

[arXiv:1809.03290]

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## ANALYSIS OVERVIEW:

- Aim: to measure  $|V_{cb}|$
- Belle sample:  $711 \text{ fb}^{-1}$
- Untagged analysis
- 3D–Binned Maximum Likelihood fit:
  - ◊  $\cos \theta_{BY}$
  - ◊  $\Delta M$ : mass difference  $D^* - D^0$
  - ◊ Lepton momentum
- Separates electron and muon channel
- Signal yields:  $\sim 90000$  for each channel

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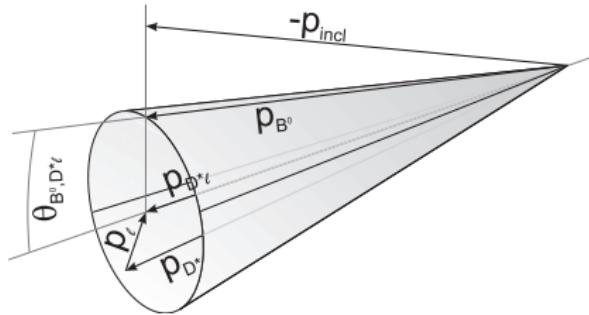
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$\theta_{BY}$ : Angle between directions of the  $B$ -meson and the  $D^* \ell (Y)$  system

$$\cos \theta_{BY} = \frac{2E_B^* E_Y^* - 2M_B^2 - m_Y^2}{2p_B^* p_Y^*} \quad (8)$$

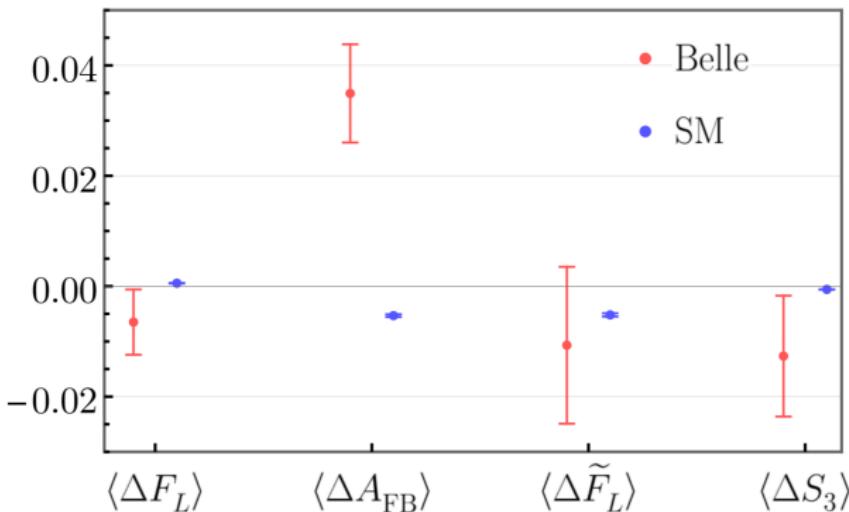


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# Belle analysis



Next step: Extraction of observables, sensitive to BSM physics

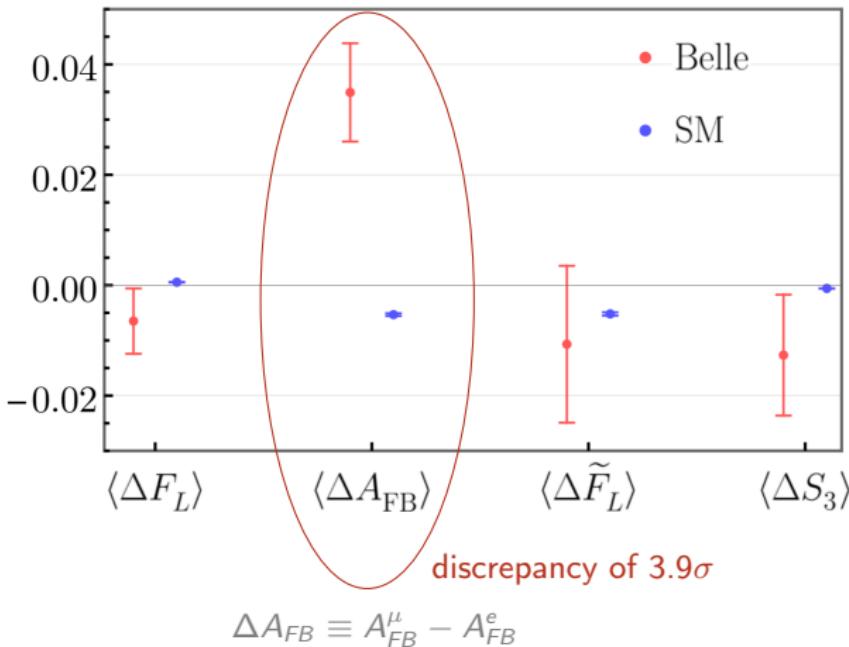


[arXiv:2104.02094v1]

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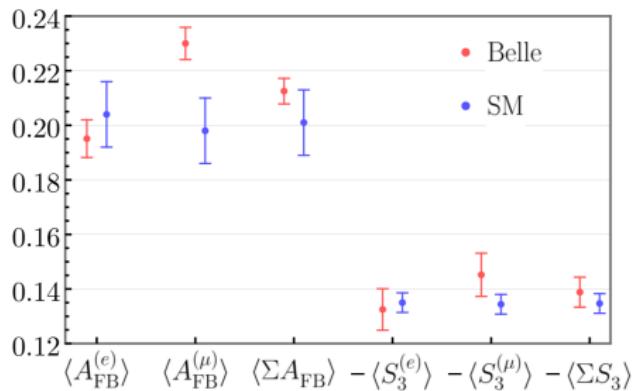
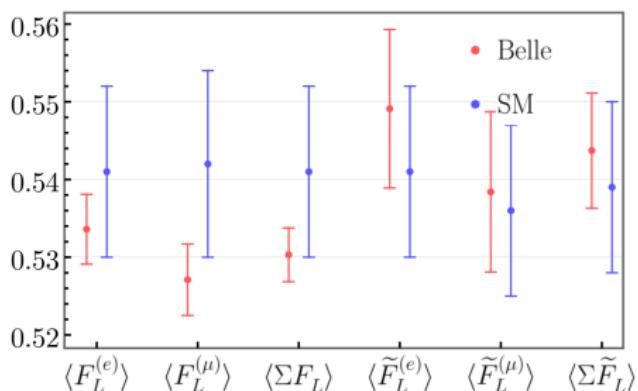


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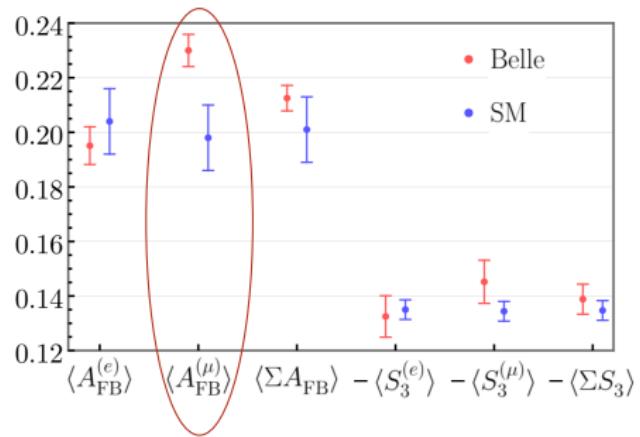
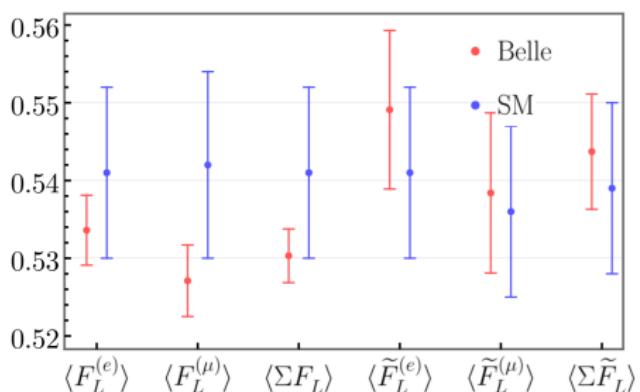


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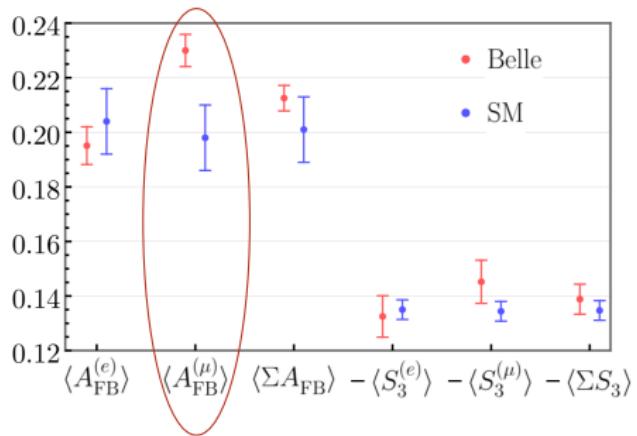
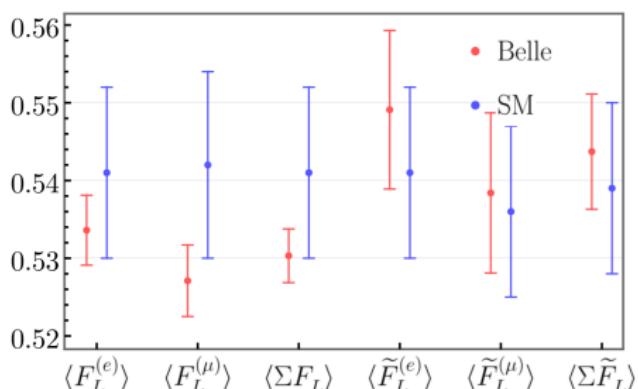
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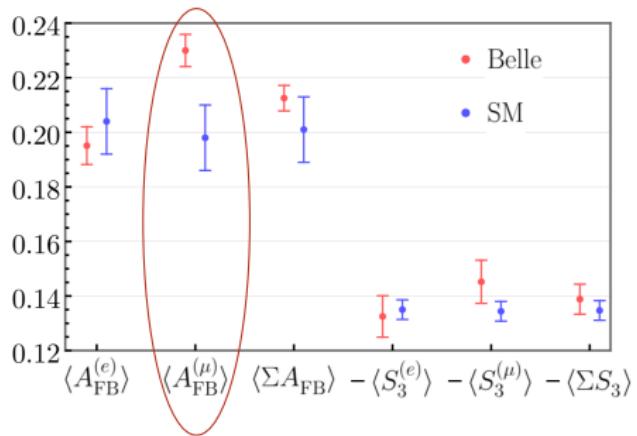
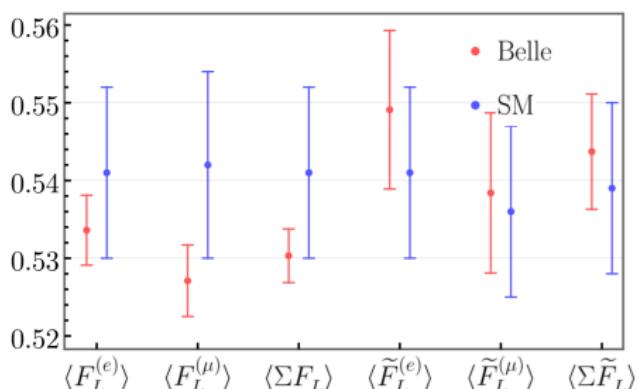


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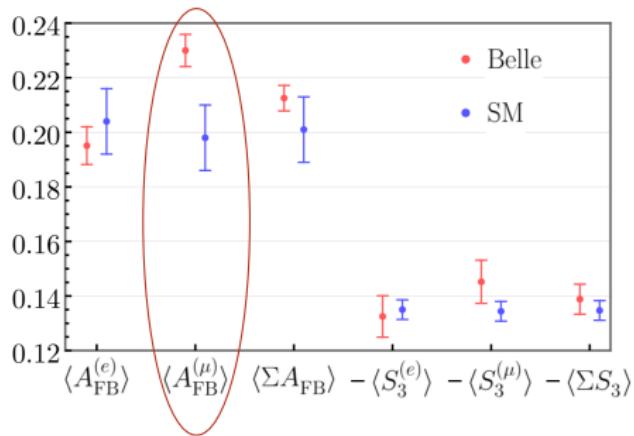
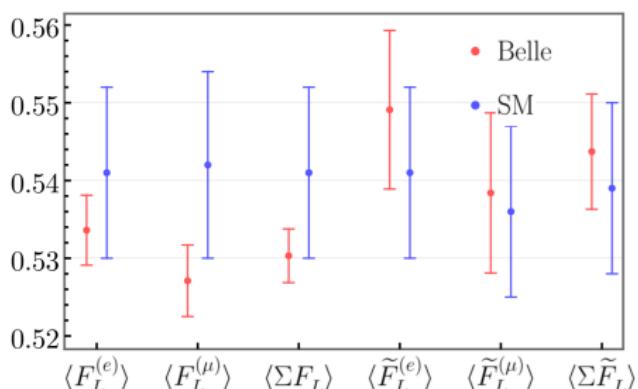
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- Correlation matrices of the statistical uncertainties are incorrect
- ? Beyond the Standard Model physics scenario
- ? Wrong assumptions → But why no discrepancy in other parameters?

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- Correlation matrices of the statistical uncertainties are incorrect
- ? Beyond the Standard Model physics scenario
- ? Wrong assumptions → But why no discrepancy in other parameters?
- Further studies are needed
- $e/\mu$  flavours should be studied separately

# Belle II Studies of $\bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$



## Main objectives

### Measurements of $|V_{cb}|$ and $A_{FB}$

Analysis overview:

- Similar to Belle analysis
- Untagged approach
- Fitting variable:  $\cos \theta_{BY}$
- Separation between  $e$  and  $\mu$  channel
- Measurements of several observables, like  $q^2$  and angle distributions

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[arXiv:2008.07198]

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<b>Event</b>	$n\text{Tracks} \geq 3$ $E_{\text{vis}}^* > 4 \text{ GeV}$ $\text{FoxWolframR2} < 0.3$
<b>Tracking</b>	$ z_0  < 2 \text{ cm}$ $ d_0  < 0.5 \text{ cm}$ $\Theta$ in CDC acceptance
<b>Brems-Photon</b>	$0.05 \leq E \leq 0.15 \text{ GeV}$
<b>Leptons</b>	$t\text{ID} > 0.9$ $1.2 < p_t^* < 2.4 \text{ GeV}$
<b>Slow pions</b>	$p_{\pi\pi}^* < 0.4 \text{ GeV}$
<b>D mesons</b>	$1.85 < m_D < 1.88 \text{ GeV}$ $0.144 < \Delta m_D < 0.148 \text{ GeV}$ $p_{D^*}^* < 2.5 \text{ GeV}$
$\cos \theta_{BY}$	$ \cos \theta_{BY}  \leq 4$
<b>For kinematic variables only</b>	$w \leq 1.5$

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## Data samples

- **Belle II data** collected in the years 2019 and 2020 equivalent to  $62.8 \text{ fb}^{-1}$
- Background and signal are modeled by **Monte Carlo sample of  $300 \text{ fb}^{-1}$**

# Untagged analysis: $B$ rest frame reconstruction



Belle II reconstruction framework uses dedicated algorithms to improve  $B$  rest frame reconstruction.

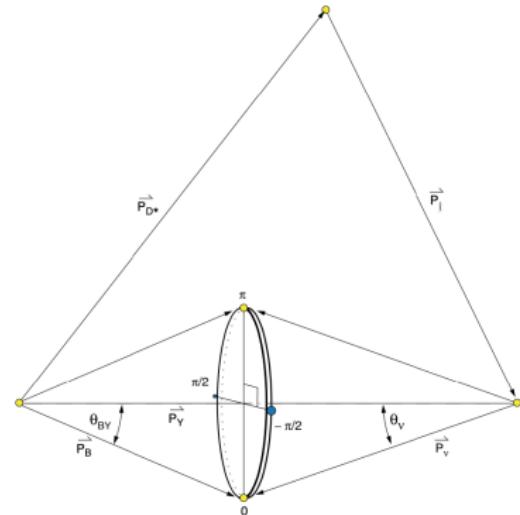
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## DIAMOND FRAME

- $B$ -mesons generated perpendicularly to the direction of the  $\Upsilon(4S)$  in  $\Upsilon(4S) \rightarrow B\bar{B}$
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- Four azimuthal angles  $\phi$  for the weighted average (compromise between computing and precision)



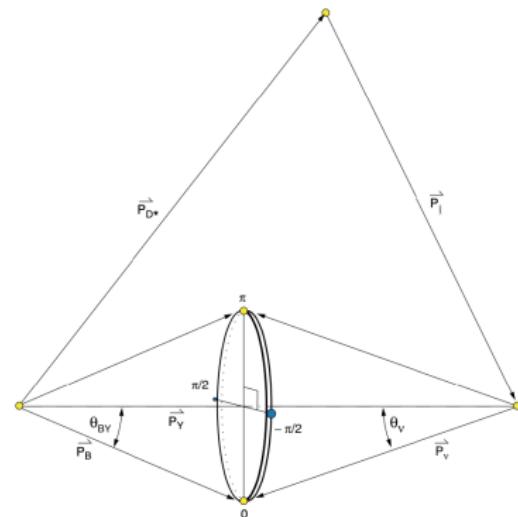
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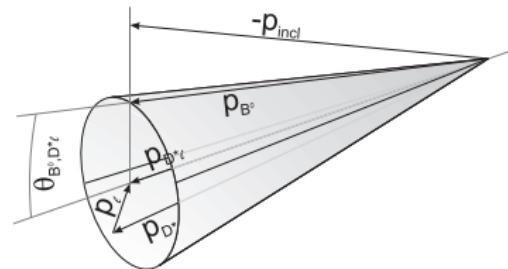
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## Rest-Of-Event (ROE) FRAME

- Estimate momentum  $\vec{p}_{\text{incl}}$  of non-signal  $B$ -meson
- Choose direction of cone based on  $\vec{p}_{\text{incl}}$

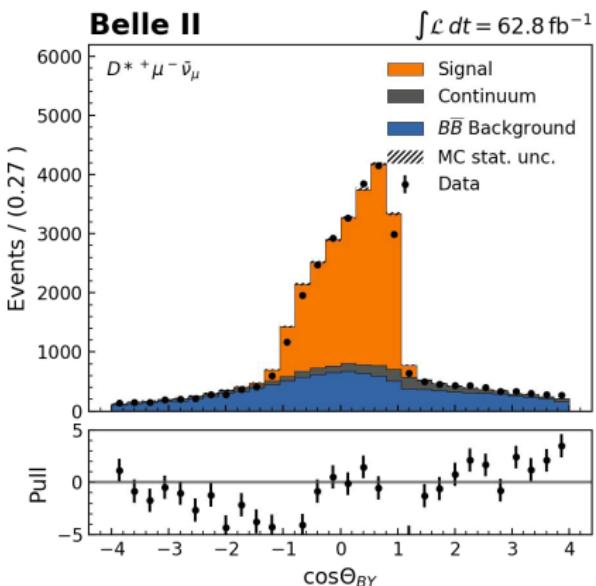
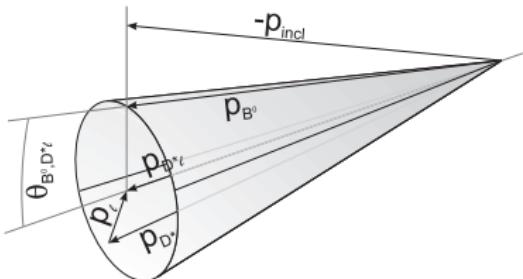


# $B^0 \rightarrow D^{*-} \ell^+ \nu_\ell$ : Fitting procedure



- Binned Likelihood Fit of  $\cos \theta_{BY}$
- Three MC templates (shapes) are used:
  - ◊ Signal events
  - ◊  $B\bar{B}$  backgrounds
  - ◊ Continuum ( $e^+ e^- \rightarrow q\bar{q}$ )

$$\cos \theta_{BY} = \frac{2E_B^* E_Y^* - 2M_B^2 - m_Y^2}{2p_B^* p_Y^*}$$



# Branching fraction



$$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell) = \frac{N_{\text{sig}}^\ell}{N_{B^0} \times \epsilon_{B^0} \times \mathcal{B}(D^{*+} \rightarrow D^0 \pi_s^+) \times \mathcal{B}(D^0 \rightarrow K^- \pi^+)} \quad (9)$$

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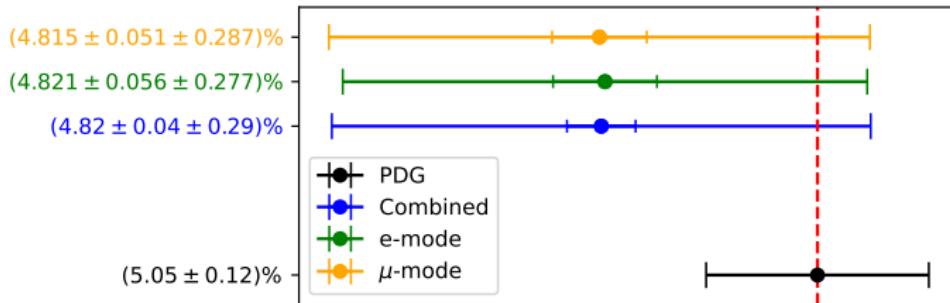
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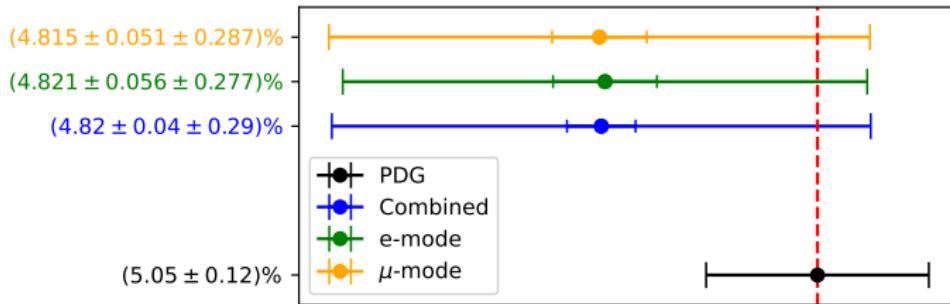
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Lepton universality:

- Probed by ratio of  $\mathcal{B}_e/\mathcal{B}_\mu = 1.001 \pm 0.016$



Back to  $A_{FB}$ ...



→ Besides the three angles  $(\cos \theta_\ell, \cos \theta_D, \chi)$  and  $q^2$ , we define one more variable,  $w$ :

$$q^2 = M_{\ell\nu}^2 = 2(E_\ell E_\nu - \vec{p}_\ell \cdot \vec{p}_\nu) \quad (10)$$

$$w = \frac{\vec{p}_B \cdot \vec{p}_{D^*}}{m_B m_{D^*}} = \frac{m_B^2 + m_{D^*}^2 - q^2}{2m_B m_{D^*}} \quad (11)$$

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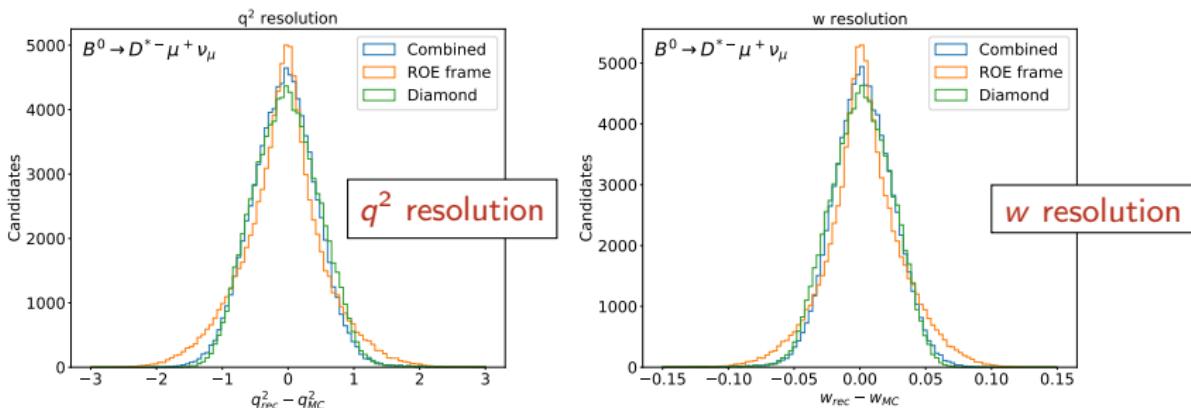
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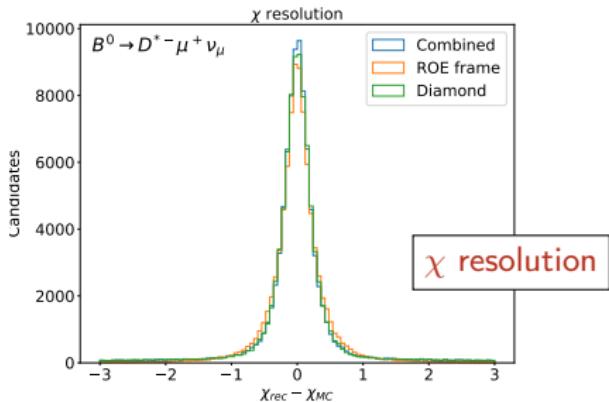
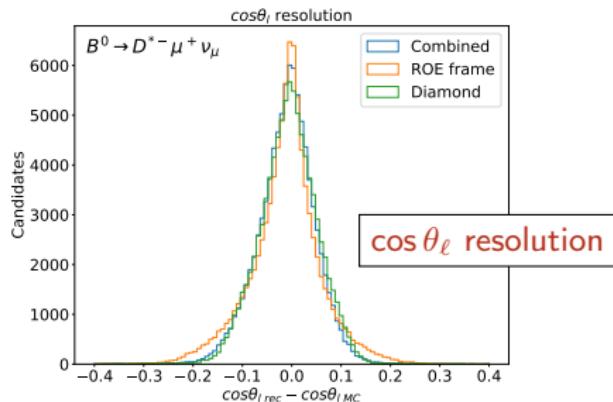
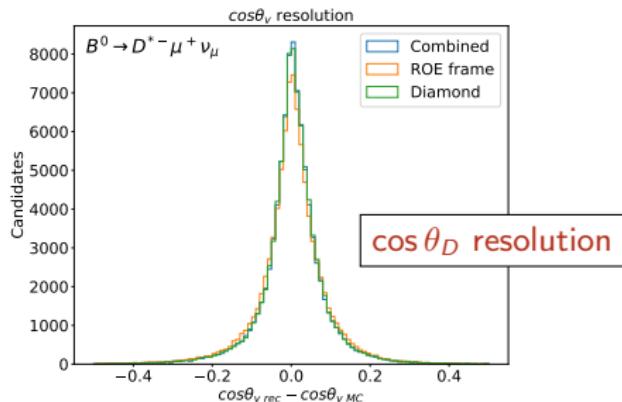
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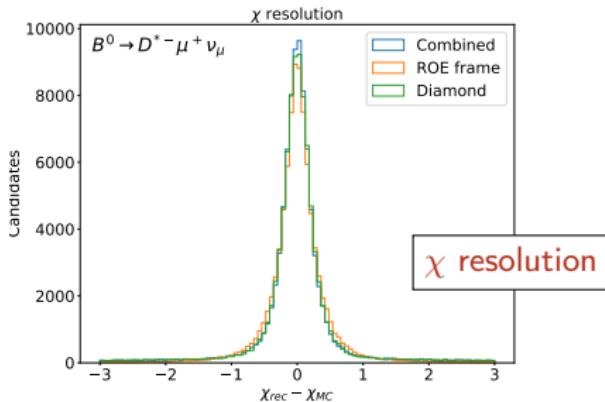
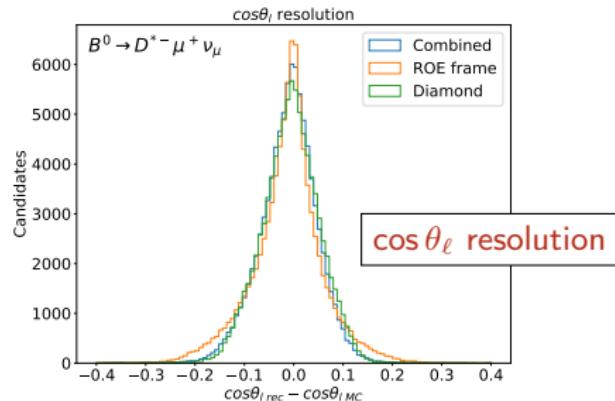
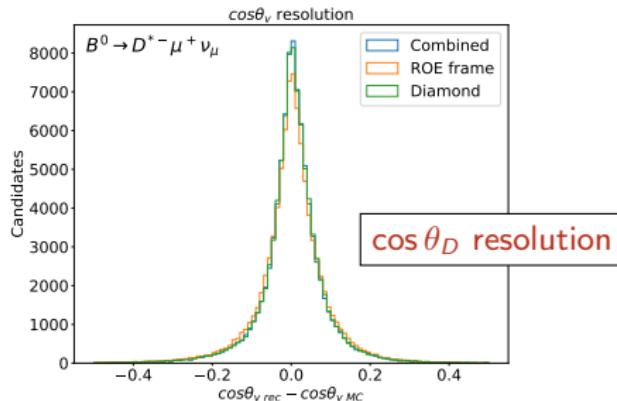
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## RESOLUTIONS

	Belle	Belle II
w	0.025	0.026
$\cos\theta_D$	0.050	0.060
$\cos\theta_\ell$	0.049	0.044
$\chi$	$13.48^\circ$ *	0.288 rad
	$* \sim 0.235$ rad	

Belle:  $141 \text{ fb}^{-1}$ , Belle II:  $300 \text{ fb}^{-1}$   
 (different methods for resolution determinations)

# Summary & Outlook



## BELLE RESULTS

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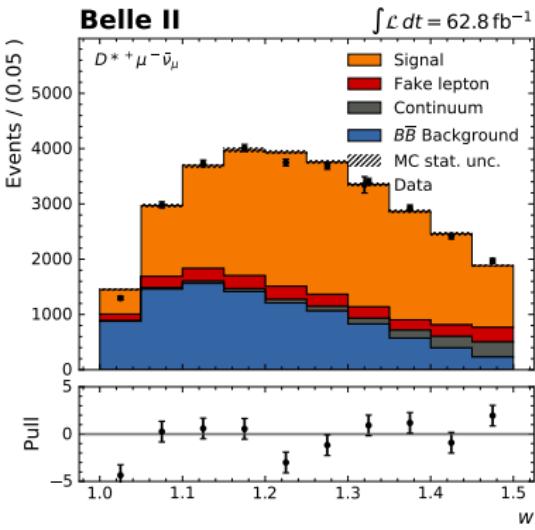
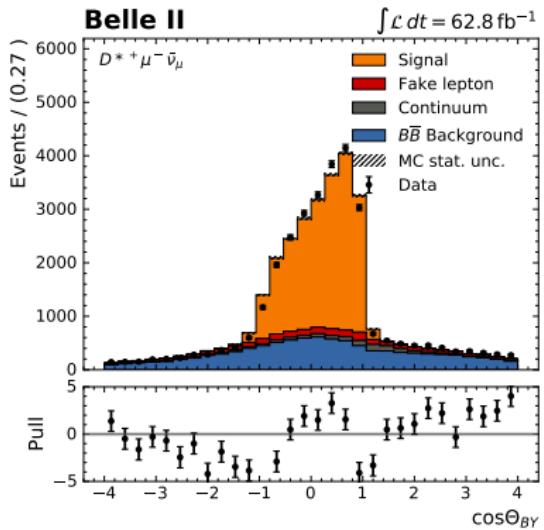
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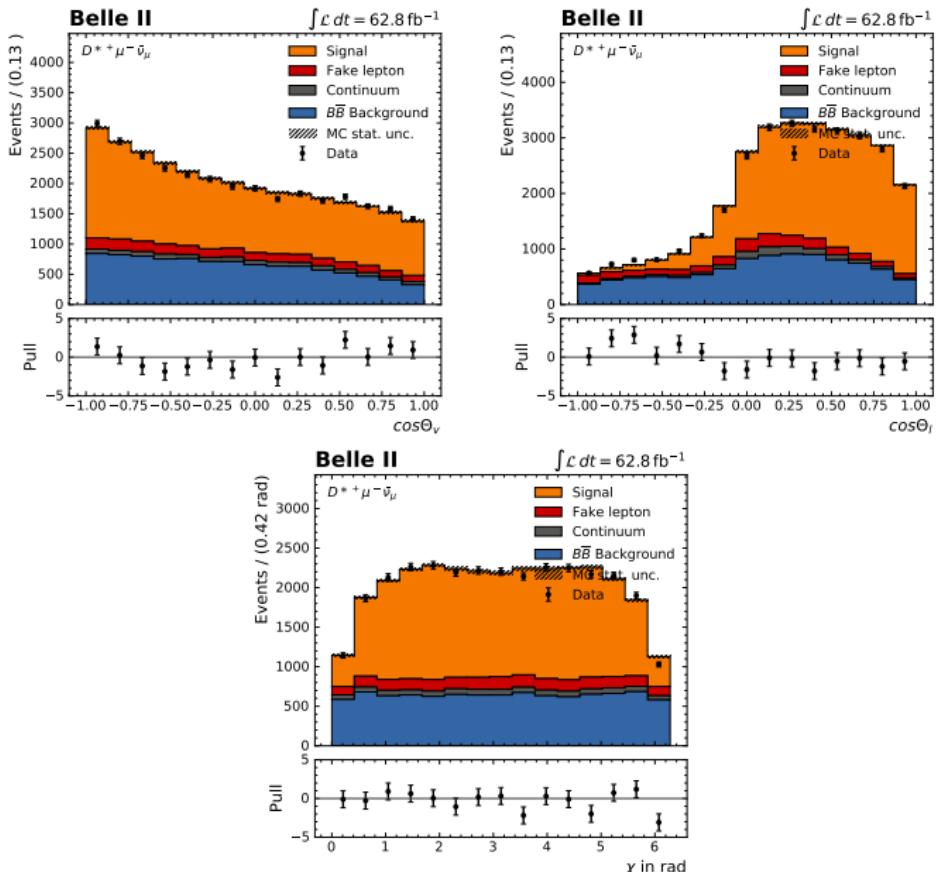
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We plan to obtain  $A_{FB}$  with Belle II data soon! :)

# BACKUP

# Distributions of observables – $\mu$ mode





# Resolution determination



Different methods were tried:

- Fit with one gaussian distributions
- Fit with two gaussian distributions
- Integrating over  $\mu \pm 1\sigma$

	Belle		Belle II	
	$\sigma_1$	$\sigma_2$	$\sigma_1$	$\sigma_2$
$\cos \theta_D$	0.033	0.086	0.026	0.080
$\cos \theta_\ell$	0.028	0.095	0.024	0.082
$\chi$	0.055	0.175	0.142	0.403

