

# The progress and prospect on charm mixing

Vitaly Vorobyev

Budker Institute of Nuclear Physics

and

Novosibirsk State University

On behalf of Belle & Belle II Collaborations

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# The charm mixing parameters

Charm forms the only neutral meson system with the heavy up quark

- Mass eigenstates

$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle, \quad |p|^2 + |q|^2 = 1$$

- The charm mixing parameters

$$x \equiv \frac{m_2 - m_1}{\Gamma}, \quad y \equiv \frac{\Gamma_2 - \Gamma_1}{2\Gamma}, \quad \Gamma \equiv \frac{\Gamma_1 + \Gamma_2}{2}$$

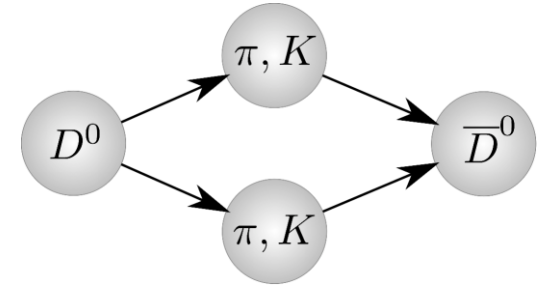
- The observable parameter

$$\lambda_f \equiv \frac{q \bar{\mathcal{A}}_f}{p \mathcal{A}_f}, \quad \mathcal{A}_f \equiv \langle f | \mathcal{H} | D^0 \rangle, \quad \bar{\mathcal{A}}_f \equiv \langle f | \mathcal{H} | \bar{D}^0 \rangle$$

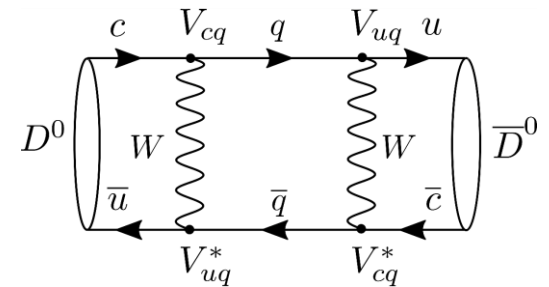
- $\mathcal{CP}$  violation

- **Direct:**  $|\langle f | \mathcal{H} | D^0 \rangle| \neq |\langle \mathcal{CP}(f) | \mathcal{H} | \bar{D}^0 \rangle|$
- **In mixing:**  $|q/p| \neq 1$
- **In interference between mixing and decay:**  $\text{Im } \lambda_f \neq 0$

Long distances  
(dominant)



Short distances



- The Standard Model expectation

$$x \lesssim y \sim \sin^2 \theta_C \times [SU(3)_f \text{ breaking}]^2$$

- Clear signals of new dynamics:

- $y \ll x \sim 1\%$
- $\mathcal{CP}$  violation  $> 10^{-3}$

# Experimental status

arXiv:1612.07233

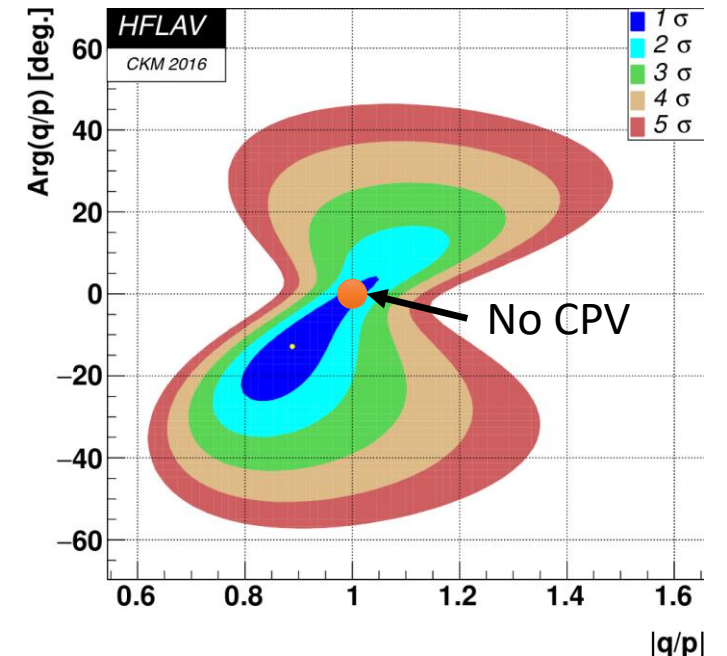
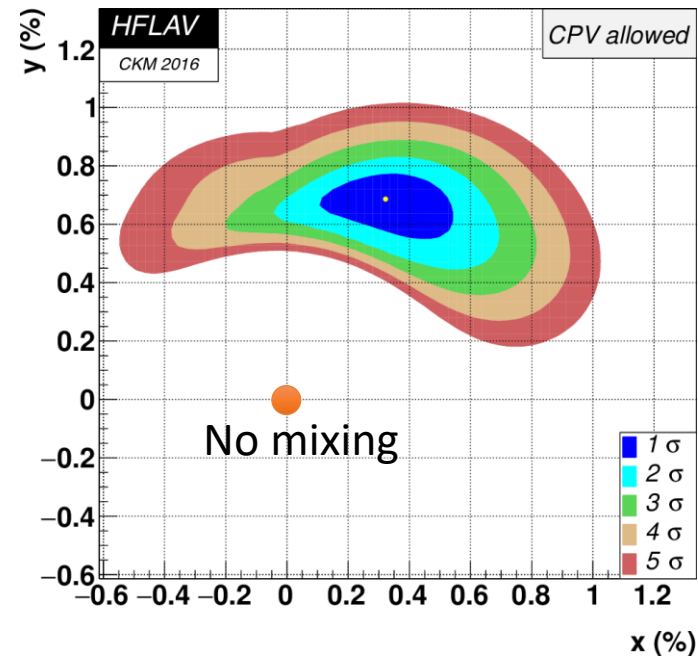
$$x = (0.32 \pm 0.14)\%$$

$$y = (0.69^{+0.06}_{-0.07})\%$$

$$|q/p| = 0.89^{+0.08}_{-0.07}$$

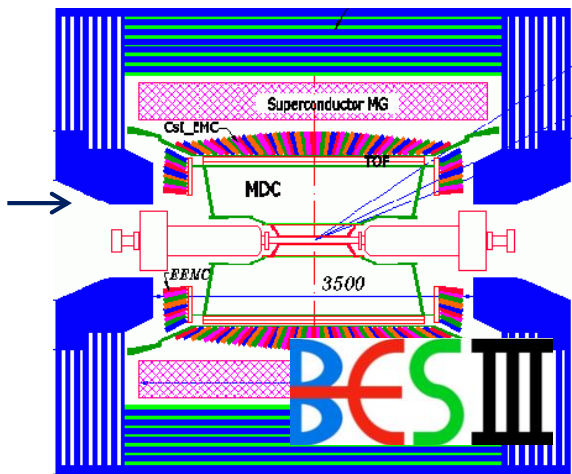
$$\arg(q/p) = (-12.9^{+9.9}_{-8.7})^\circ$$

- The first evidence by Belle and BaBar
  - Phys. Rev. Lett. 98, 211802 (2007)
  - Phys. Rev. Lett. 98, 211803 (2007)
- Charm mixing is well established
- No  $\mathcal{CP}$  violation observed in charm yet

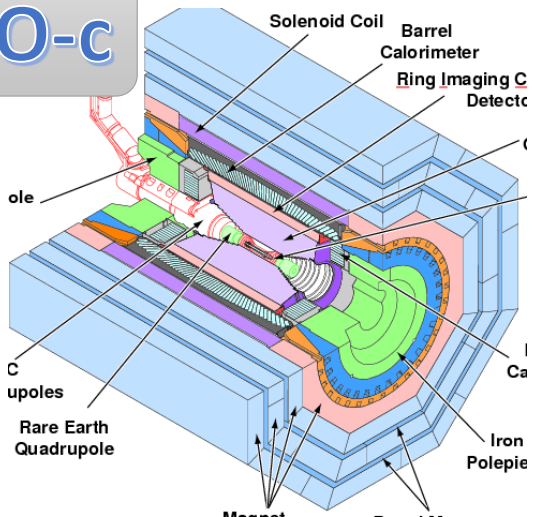


# Experimental landscape

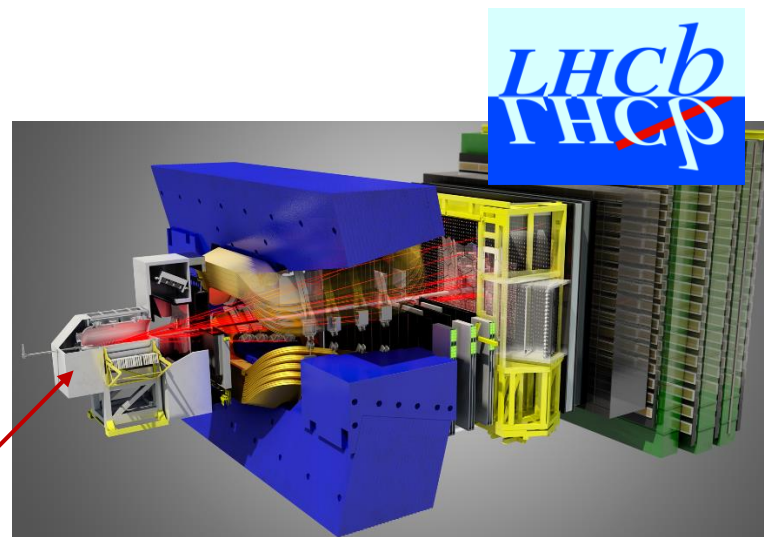
$D\bar{D}$  threshold



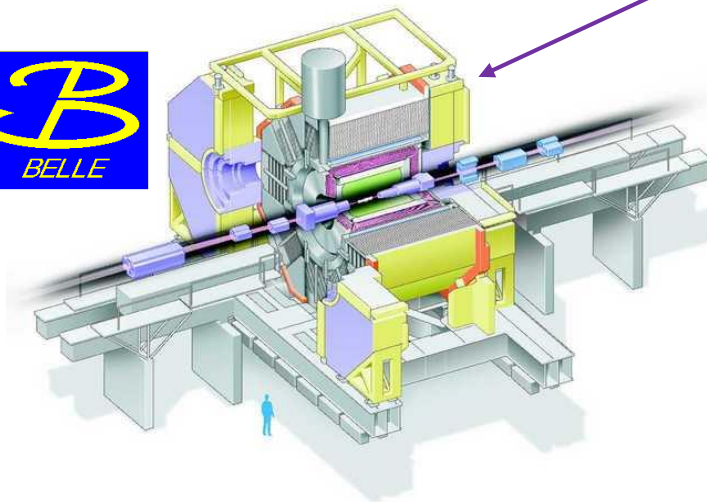
CLEO-c



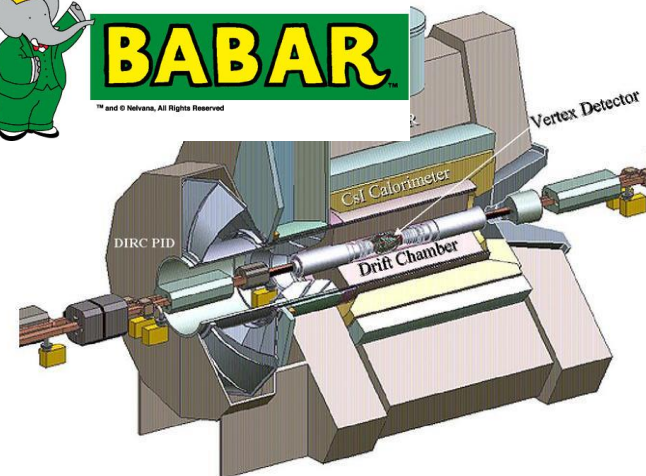
$pp$  collisions



Asymmetric  $e^+e^-$  @  $\Upsilon(4S)$



**BABAR**



# Charm mixing observables and facilities

## Classes of observables (I. Bigi)

'Wrong-sign'  $D^0$  decays

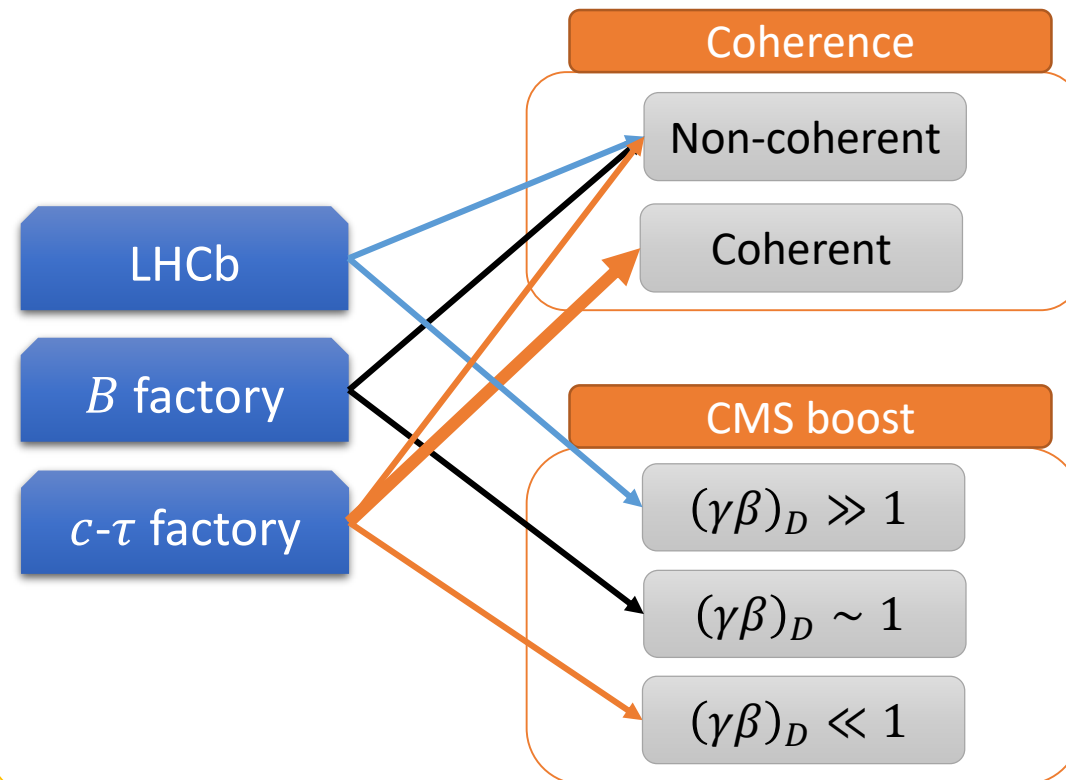
- Semileptonic
- Hadronic ( $D^0 \rightarrow K^+\pi^-$  time ind.)

Non-exponential  $D^0$  decay rate evolutions ( $D^0 \rightarrow K^+\pi^-$  time dep.)

Different  $D^0$  lifetimes in different channels ( $D^0 \rightarrow h^+h^-$ )

$D^0$  and  $\bar{D}^0$  produced in  $e^+e^- \rightarrow D^0\bar{D}^0$  decaying into two seemingly identical final states

## Charm production



# Charm decay rates

Time-dependent

Incoherent

$$D^{*\pm} \rightarrow D\pi^\pm, \quad B \rightarrow DX, \quad e^+e^- \rightarrow c\bar{c} \rightarrow D\bar{D}X, \quad pp \rightarrow c\bar{c}X$$

$$|\langle f | \mathcal{H} | D^0(t) \rangle|^2 = e^{-\Gamma t} |\mathcal{A}_f|^2 [1 - (y \operatorname{Re}\lambda_f + x \operatorname{Im}\lambda_f)\Gamma t] + \mathcal{O}(x^2, y^2)$$

$$|\langle f | \mathcal{H} | D^0 \rangle|^2 \propto |\mathcal{A}_f|^2 (1 - y \operatorname{Re}\lambda_f - x \operatorname{Im}\lambda_f) + \mathcal{O}(x^2, y^2)$$

Boost

$$\text{LHCb: } (\gamma\beta)_D \gg 1$$

$$\text{B factory: } (\gamma\beta)_D \sim 1$$

$$\text{c-}\tau \text{ factory: } (\gamma\beta)_D \ll 1$$

Time-integrated

Coherent (at rest)

$$e^+e^- \rightarrow D^{(*)0}\bar{D}^{(*)0}, \quad \mathcal{C}+: D^0\bar{D}^0\gamma, \quad \mathcal{C}-: D^0\bar{D}^0(\pi^0)$$

$$\langle ij | \mathcal{H} | D^0\bar{D}^0 \rangle \propto \langle i | \mathcal{H} | D^0 \rangle \langle j | \mathcal{H} | \bar{D}^0 \rangle + \mathcal{C} \langle i | \mathcal{H} | \bar{D}^0 \rangle \langle j | \mathcal{H} | D^0 \rangle$$

$$|\langle ij | \mathcal{H} | D^0\bar{D}^0 \rangle|^2 \propto |\mathcal{A}_i|^2 |\mathcal{A}_j|^2 [|\zeta_c|^2 + (1 + \mathcal{C})(x \operatorname{Im}(\xi_c^* \zeta_c) - y \operatorname{Re}(\xi_c^* \zeta_c))] + \mathcal{O}(x^2, y^2)$$

$$\xi_c \equiv \frac{p}{q}(1 + \mathcal{C}\lambda_i\lambda_j), \quad \zeta_c \equiv \frac{p}{q}(\lambda_j + \mathcal{C}\lambda_i)$$

# The progress

Selected experimental results

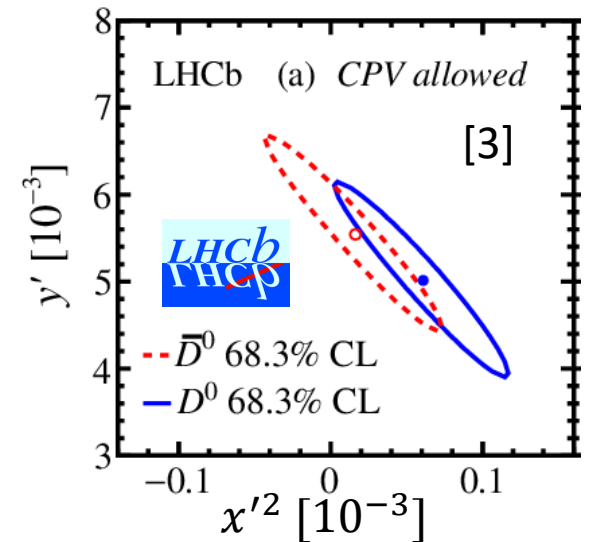
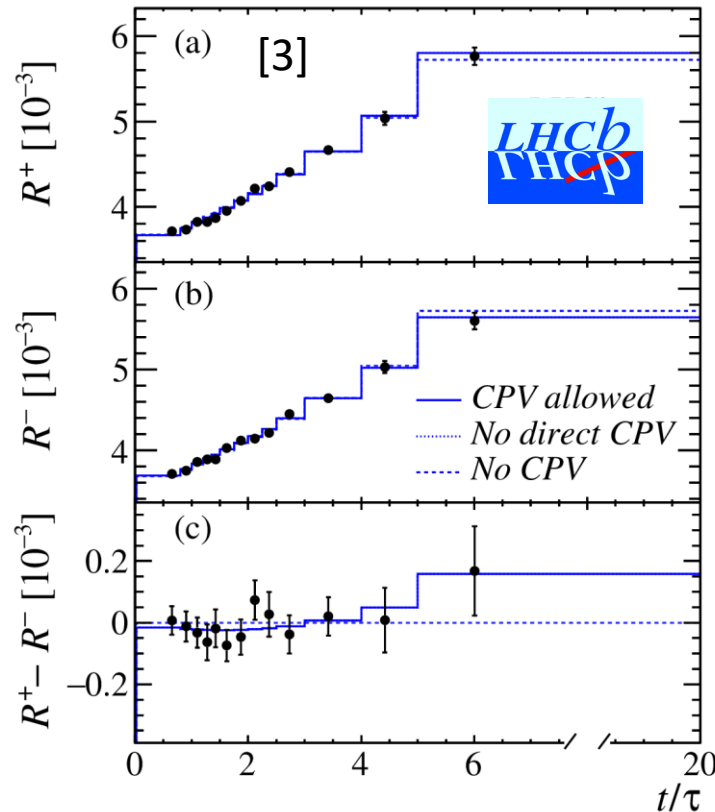
# $D^0 \rightarrow K\pi$ time-dependent WS

$$\Gamma(D^0(t) \rightarrow f_{WS}) = e^{-\frac{t}{\tau}} |A_f|^2 \left[ R_D + \sqrt{R_D} y' \frac{t}{\tau} + \frac{1}{2} R_M \left( \frac{t}{\tau} \right)^2 \right]$$

$$\begin{aligned} x' &\equiv x \cos \delta_{K\pi} + y \sin \delta_{K\pi} \\ y' &\equiv y \cos \delta_{K\pi} - x \sin \delta_{K\pi} \\ R_M &\equiv (x^2 + y^2)/2 \end{aligned}$$

$$\begin{aligned} R_D^{(+)} &\equiv \frac{Br(D^0 \rightarrow K^+ \pi^-)}{Br(D^0 \rightarrow K^- \pi^+)} \\ A_D &\equiv \frac{(R_D^+ - R_D^-)}{(R_D^+ + R_D^-)} \end{aligned}$$

- BaBar [1]:  $384 \text{ fb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$   
 $y' = (9.7 \pm 4.4 \pm 3.1) \times 10^{-3}$   
 $R_D = (3.03 \pm 0.19) \times 10^{-3}$
- Belle [2]:  $976 \text{ fb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$   
 $y' = (4.6 \pm 3.4) \times 10^{-3}$   
 $R_D = (3.53 \pm 0.13) \times 10^{-3}$
- LHCb [3]:  $5 \text{ fb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$   
 $y' = (5.28 \pm 0.45 \pm 0.27) \times 10^{-3}$   
 $x'^2 = (0.039 \pm 0.023 \pm 0.014) \times 10^{-3}$   
 $R_D = (3.454 \pm 0.028 \pm 0.014) \times 10^{-3}$   
 $A_D = (-0.1 \pm 9.1) \times 10^{-3}$   
 $1.00 < |p/q| < 1.35 @ 68.3\% \text{ CL}$



- [1] Phys. Rev. Lett. 98, 211802 (2007) (BaBar)  
 [2] Phys. Rev. Lett. 112, 111801 (2014) (Belle)  
 [3] Phys. Rev. D97, 031101(R) (2018) (LHCb)



# Measurements with quantum correlations

Time-integrated

The method exploits the difference between correlated (double tagged) and uncorrelated (single tagged) decay rates

$\mathcal{C} = -1$  correlations

$$\Gamma(i, j) \propto |\langle i|D_2\rangle\langle j|D_1\rangle - \langle i|D_1\rangle\langle j|D_2\rangle|^2 + \mathcal{O}(x^2, y^2)$$

- CLEO-c [1]:  $0.82 \text{ fb}^{-1}$  @  $\psi(3770)$ , fit of 261 yields

$$y = (4.2 \pm 2.0 \pm 1.0)\%$$

$$R_D = (0.533 \pm 0.107 \pm 0.045)\%$$

$$\cos \delta_{K\pi} = +0.81 \pm 0.22 \pm 0.07$$

$$\sin \delta_{K\pi} = -0.01 \pm 0.41 \pm 0.04$$

- BESIII:  $2.92 \text{ fb}^{-1}$  @  $\psi(3770)$

- $y$  and  $R_D$  are taken as an external input

$$\cos \delta_{K\pi} = 1.02 \pm 0.11 \pm 0.06 \pm 0.01$$

TABLE III.  $D$  final states reconstructed in this analysis. [1]

Type	Reconstruction	Final states
$f$	Full	$K^- \pi^+, Y_0 - Y_7$
$\bar{f}$	Full	$K^+ \pi^-, \bar{Y}_0 - \bar{Y}_7$
$S_+$	Full	$K^+ K^-, \pi^+ \pi^-, K_S^0 \pi^0 \pi^0$
$S_+$	Partial	$K_L^0 \pi^0, K_L^0 \eta, K_L^0 \omega$
$S_-$	Full	$K_S^0 \pi^0, K_S^0 \eta, K_S^0 \omega$
$S_-$	Partial	$K_L^0 \pi^0 \pi^0$
$\ell^+$	Partial	$K^- e^+ \nu_e, K^- \mu^+ \nu_\mu$
$\ell^-$	Partial	$K^+ e^- \bar{\nu}_e, K^+ \mu^- \bar{\nu}_\mu$

[1] Phys. Rev. D86, 112001 (2012) (CLEO-c)

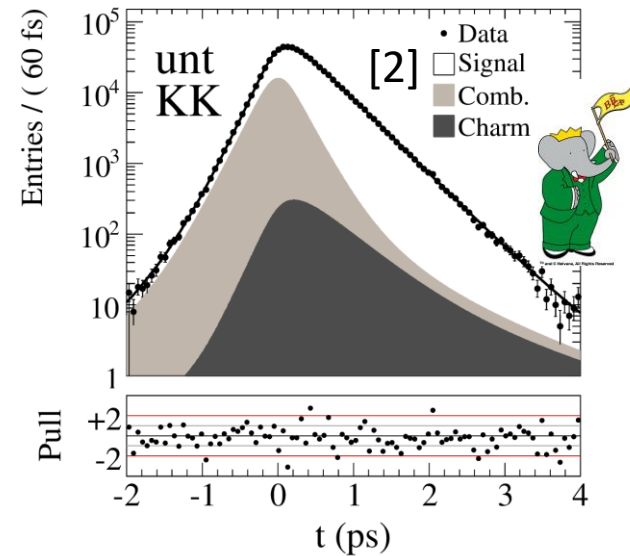
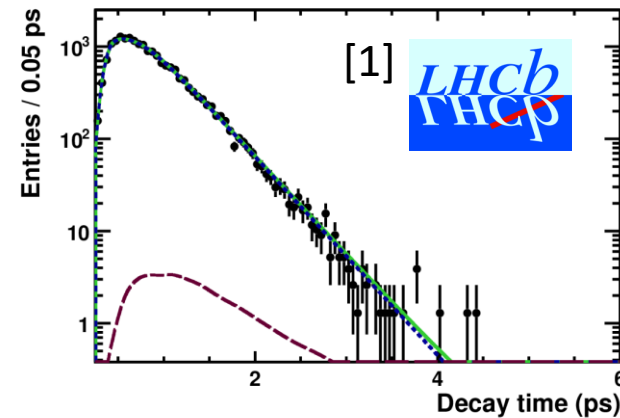
[2] Phys. Lett. B 734, 277 (2014) (BESIII)

# Time-dependent $D^0 \rightarrow h^+ h^-$

$$y_{CP} \equiv \eta_{CP} \frac{\hat{\Gamma}(D^0 \rightarrow f) + \hat{\Gamma}(\bar{D}^0 \rightarrow f)}{2\hat{\Gamma}(D^0 \rightarrow K^- \pi^+)} - 1 \approx y \cos \varphi - \frac{1}{2} A_m x \sin \varphi$$

$$y_{CP} \approx \frac{\tau(D^0 \rightarrow K^- \pi^+)}{\tau(D^0 \rightarrow K^- K^+)} - 1$$

- LHCb [1]\*:  $29 \text{ pb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$   
 $y_{CP} = (0.55 \pm 0.63 \pm 0.41)\%$
- BaBar [2]:  $468 \text{ fb}^{-1}$ ,  $D^0 \rightarrow K^{\mp} \pi^{\pm}, K^- K^+, \pi^- \pi^+$   
 $y_{CP} = (0.72 \pm 0.18 \pm 0.12)\%$
- Belle [3]:  $540 \text{ fb}^{-1}$ ,  $D^0 \rightarrow K^{\mp} \pi^{\pm}, K^- K^+, \pi^- \pi^+$   
 $y_{CP} = (1.31 \pm 0.32 \pm 0.25)\%$



[1] JHEP 04, 129 (2012) (LHCb)

[2] Phys. Rev. D87, 012004 (2013) (BaBar)

[3] Phys. Rev. Lett. 98, 211803 (2007) (Belle)

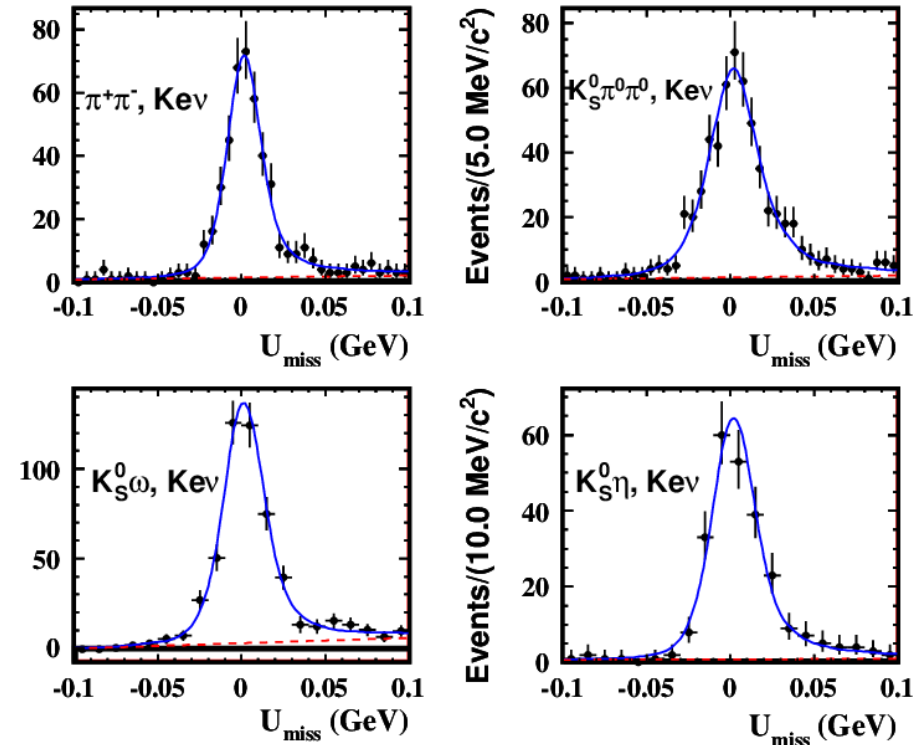
\*There are many other LHCb publications on  $CP$  violation in the  $D^0 \rightarrow h^+ h^-$  modes

# $D^0 \rightarrow f_{CP}$ with quantum correlations

$$y_{CP} \approx \frac{1}{4} \left( \frac{\mathcal{B}(D_{CP-} \rightarrow l)}{\mathcal{B}(D_{CP+} \rightarrow l)} - \frac{\mathcal{B}(D_{CP+} \rightarrow l)}{\mathcal{B}(D_{CP-} \rightarrow l)} \right), \quad \mathcal{B}(D_{CP\mp} \rightarrow l) = \frac{N_{CP\pm;l}}{N_{CP\pm}} \cdot \frac{\varepsilon_{CP\pm}}{\varepsilon_{CP\pm;l}}$$



- BESIII [1]:  $2.92 \text{ fb}^{-1}$  @  $3.773 \text{ GeV}$   
 $y_{CP} = (-2.0 \pm 1.3 \pm 0.7)\%$ 
  - Single tag:  $D \rightarrow f_{CP}$
  - Quantum correlated  
 $D\bar{D} \rightarrow f_{CP} + Kl\mu$
  - Systematic uncertainty has statistical origin



[1] Phys. Lett. B744, 339 (2015)

# $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ Dalitz plot analysis

- Time-dependent Dalitz plot analysis

$$\mathcal{P}_D(t, m_+^2, m_-^2) \approx \Gamma e^{-\Gamma t} [|\mathcal{A}_D|^2 - \Gamma t \operatorname{Re}(\mathcal{A}_D^* \mathcal{A}_{\bar{D}}(y + ix))]$$

- Sensitivity to the charm mixing parameters due to the strong phase variation over the Dalitz plot

- $\mathcal{A}(m_+^2, m_-^2)$  from a  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  **decay model**

- BaBar [1]:  $468.5 \text{ fb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$

$$x = (0.16 \pm 0.23 \pm 0.12 \pm 0.08)\%$$

$$y = (0.57 \pm 0.20 \pm 0.13 \pm 0.07)\%$$

- Belle [2]:  $921 \text{ fb}^{-1}$ ,  $D^{*+} \rightarrow D^0 \pi^+$

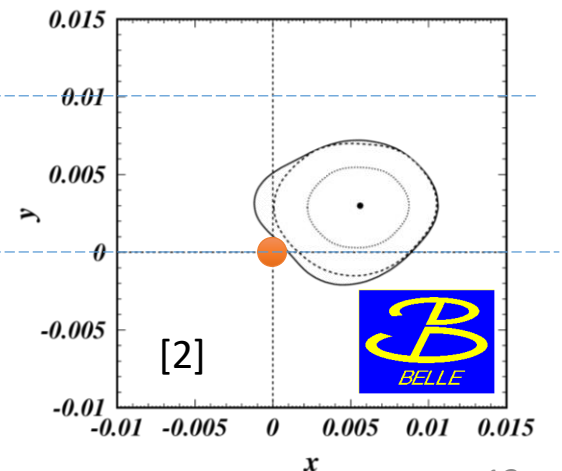
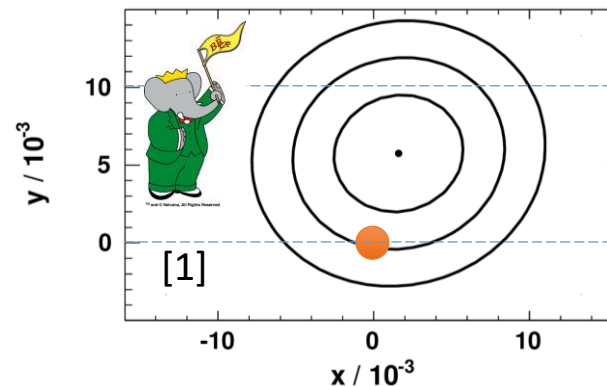
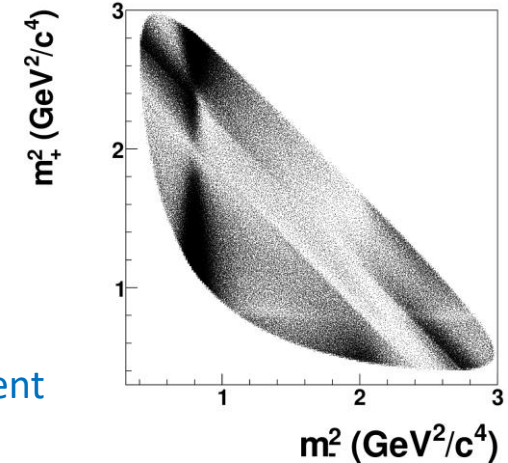
$$x = (0.56 \pm 0.19 \pm 0.08 \pm 0.08)\%$$

$$y = (0.30 \pm 0.15 \pm 0.05 \pm 0.07)\%$$

$$|q/p| = 0.90 \pm 0.16 \pm 0.05 \pm 0.06$$

$$\arg(q/p) = (-6 \pm 11 \pm 3 \pm 4)^\circ$$

The most precise single measurement of the charm mixing parameters



[1] Phys. Rev. Lett. 105, 081803 (2010) (BaBar)

[2] Phys. Rev. D89, 091103 (2014) (Belle)

# Model-independent Dalitz plot analysis

A way to eliminate (difficult to control) model dependency of a multibody decay analysis

- Binned time-dependent Dalitz plot analysis [1,2]
 
$$\mathcal{P}_D(t, i) \propto e^{-\Gamma t} [K_i - \Gamma t \sqrt{K_i K_{-i}} (C_i y + S_i x)]$$

$$\mathcal{P}_{\bar{D}}(t, i) \propto e^{-\Gamma t} [K_{-i} - \Gamma t \sqrt{K_i K_{-i}} (C_i y - S_i x)]$$
- $C_i$  and  $S_i$  are measured in coherent  $D^0 \bar{D}^0$  pair decays [3]

- LHCb [4]:  $1.0 \text{ fb}^{-1}$  @  $7 \text{ TeV}$ ,  $D^{*+} \rightarrow D^0 \pi^+$ 

$$x = (-0.86 \pm 0.53 \pm 0.17)\%$$

$$y = (+0.03 \pm 0.46 \pm 0.13)\%$$



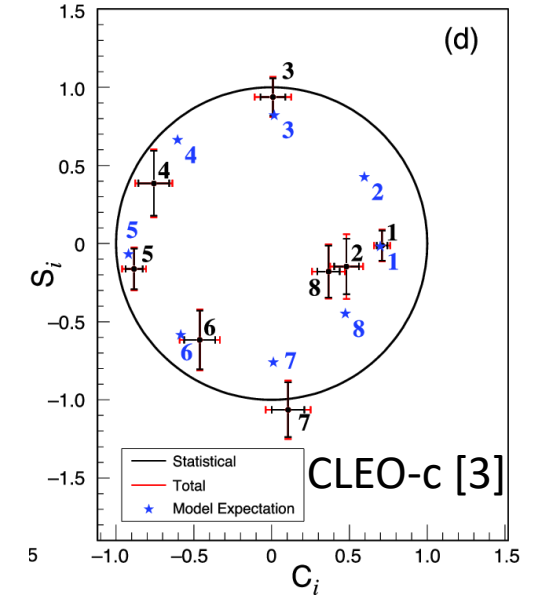
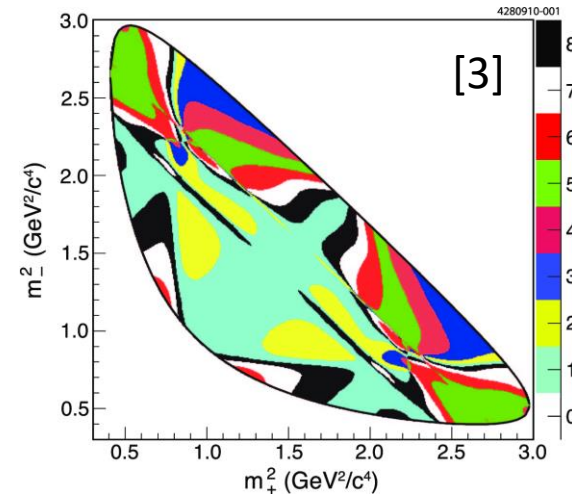
[1] Phys. Rev. D68, 054018 (2003)

[2] Phys. Rev. D82, 034033 (2010) (Bondar et al.)

[3] Phys. Rev. D82, 112006 (2010) (CLEO-c)

[4] JHEP 04, 033 (2016) (LHCb)

The first model-independent measurement of the charm mixing parameters



$$Z_i = \frac{\int_{D_i} \mathcal{A}_D^* \mathcal{A}_{\bar{D}} dm_+^2 dm_-^2}{\sqrt{\int_{D_i} |\mathcal{A}_D|^2 dm_+^2 dm_-^2 \cdot \int_{D_i} |\mathcal{A}_{\bar{D}}|^2 dm_+^2 dm_-^2}}$$

$$C_i = \text{Re } Z_i, \quad S_i = \text{Im } Z_i$$

# $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ with quantum correlations

Coherent  $\mathcal{C} = \pm 1$  and non-coherent decays

$$e^+e^- \rightarrow \psi(4040) \rightarrow D\bar{D}^*$$

- Coherent  $\mathcal{C} = -1$ :  $D^0\bar{D}^{*0} \rightarrow D^0\bar{D}^0\pi^0$

$$M_{ij}^- = K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (C_i C_j + S_i S_j)$$

- Coherent  $\mathcal{C} = +1$ :  $D^0\bar{D}^{*0} \rightarrow D^0\bar{D}^0\gamma$

$$M_{ij}^+ = K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (C_i C_j + S_i S_j) + 2K_j \sqrt{K_i K_{-i}} (yC_i - xS_i) + 2K_{-j} \sqrt{K_i K_{-i}} (yC_i + xS_i) + 2K_i \sqrt{K_j K_{-j}} (yC_j - xS_j) + 2K_{-i} \sqrt{K_j K_{-j}} (yC_j + xS_j)$$

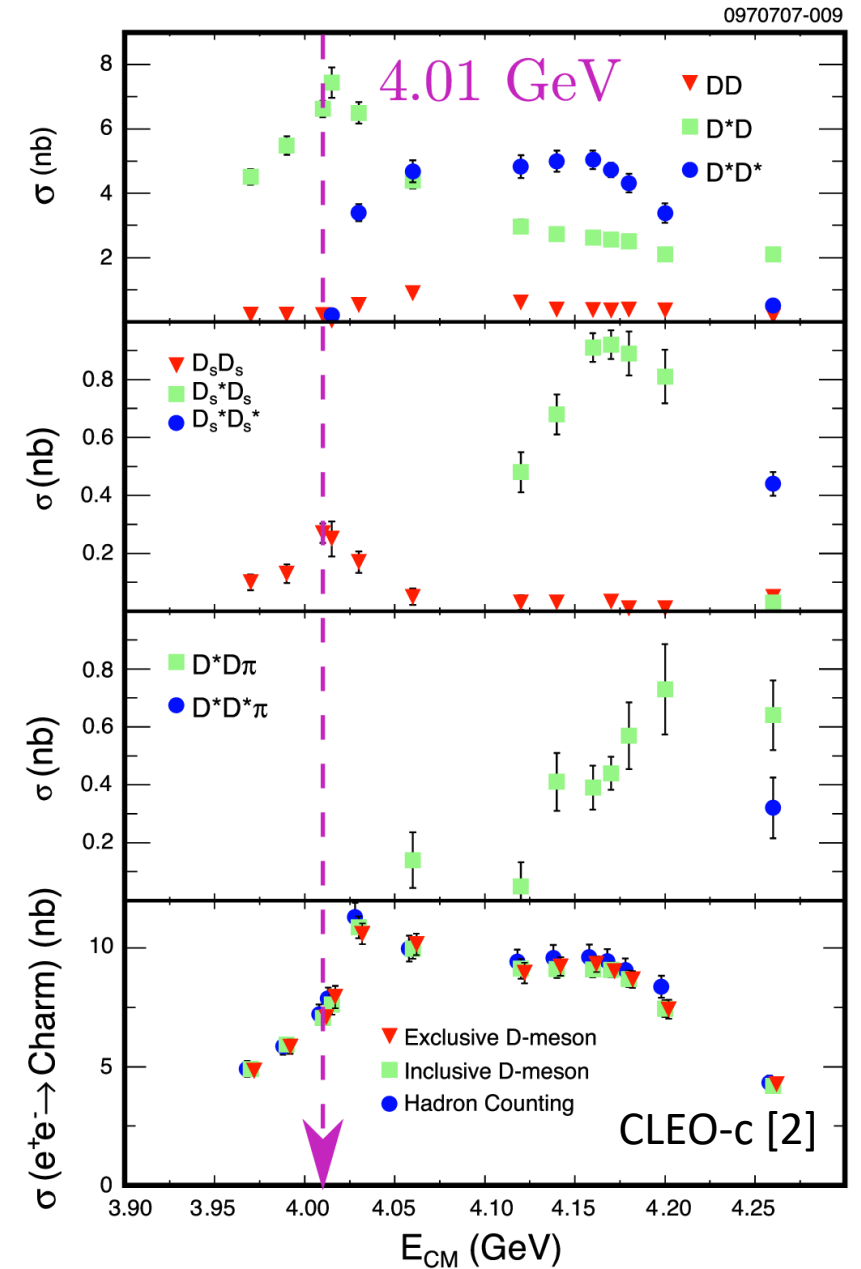
- Incoherent  $D^-D^{*+} \rightarrow D^-D^0\pi^+$

$$K_i' = K_i + \sqrt{K_i K_{-i}} (yC_i + xS_i)$$

Measurement of the charm mixing and the phase parameters in a single experiment

[1] Phys. Rev. D82, 034033 (2010)

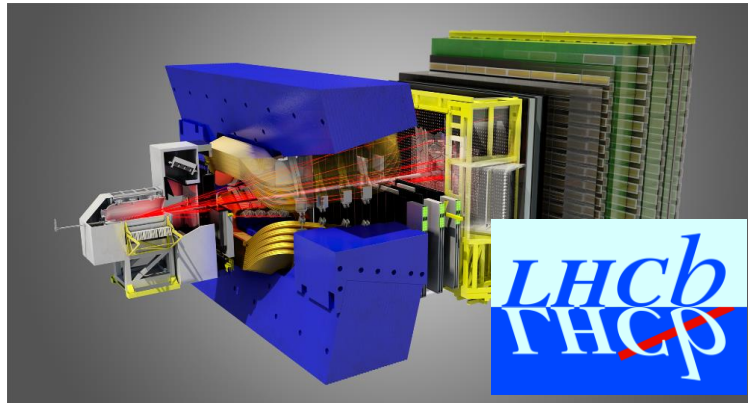
[2] Phys. Rev. D80, 072001 (2009) (CLEO-c)



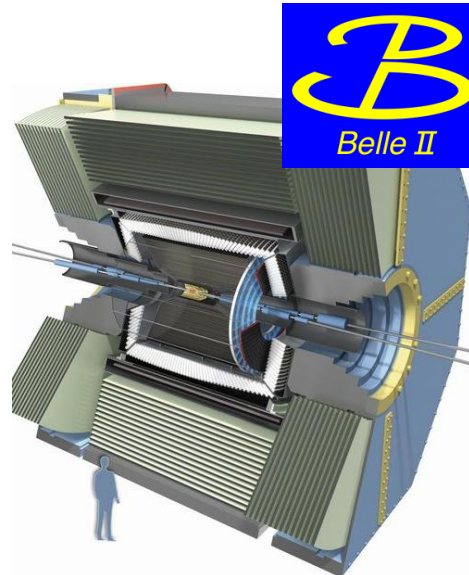
# The prospects

Estimates and expectations

# Future landscape

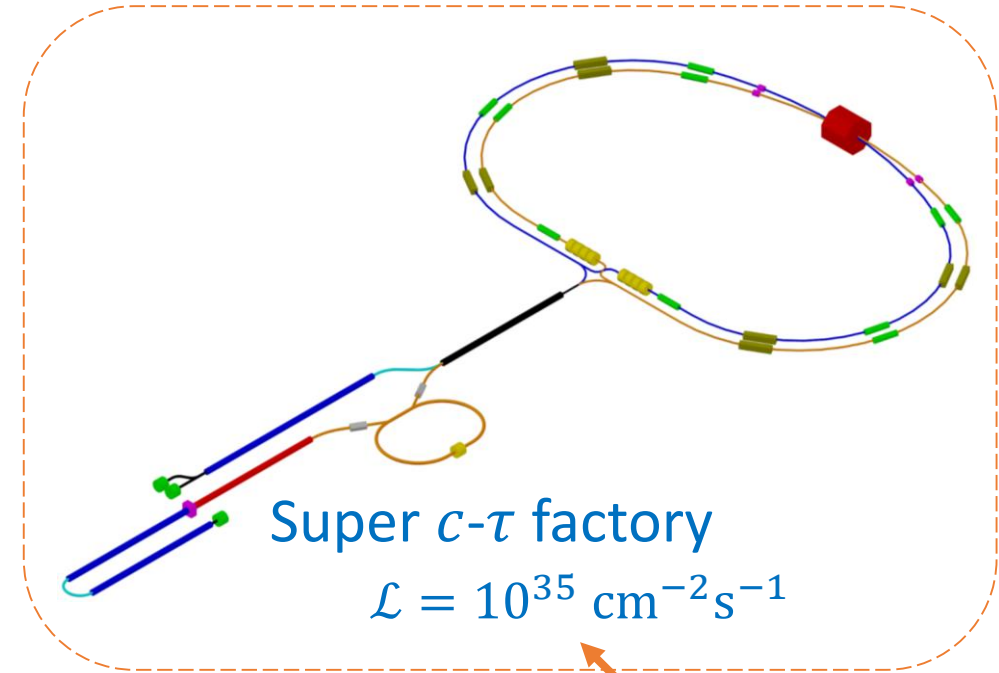


- Detector upgrade
- Trigger improvements
- Operation at higher luminosity (tens of  $pp$  interactions per event)



- Increased hermeticity and  $K_S^0$  efficiency
- Improved IP and  $D$  vertex resolution,  $K/\pi$  separation and  $\pi^0$  reconstruction
- Added PID and  $\mu$  ID in end caps

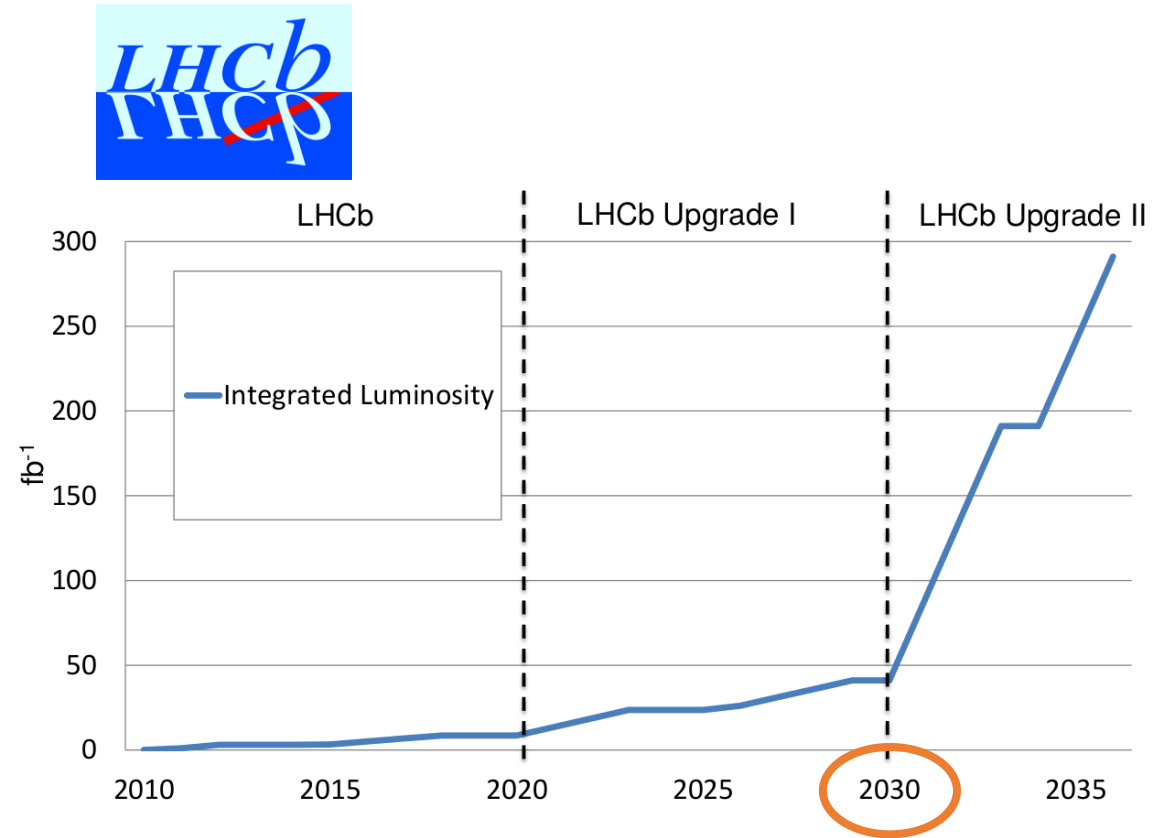
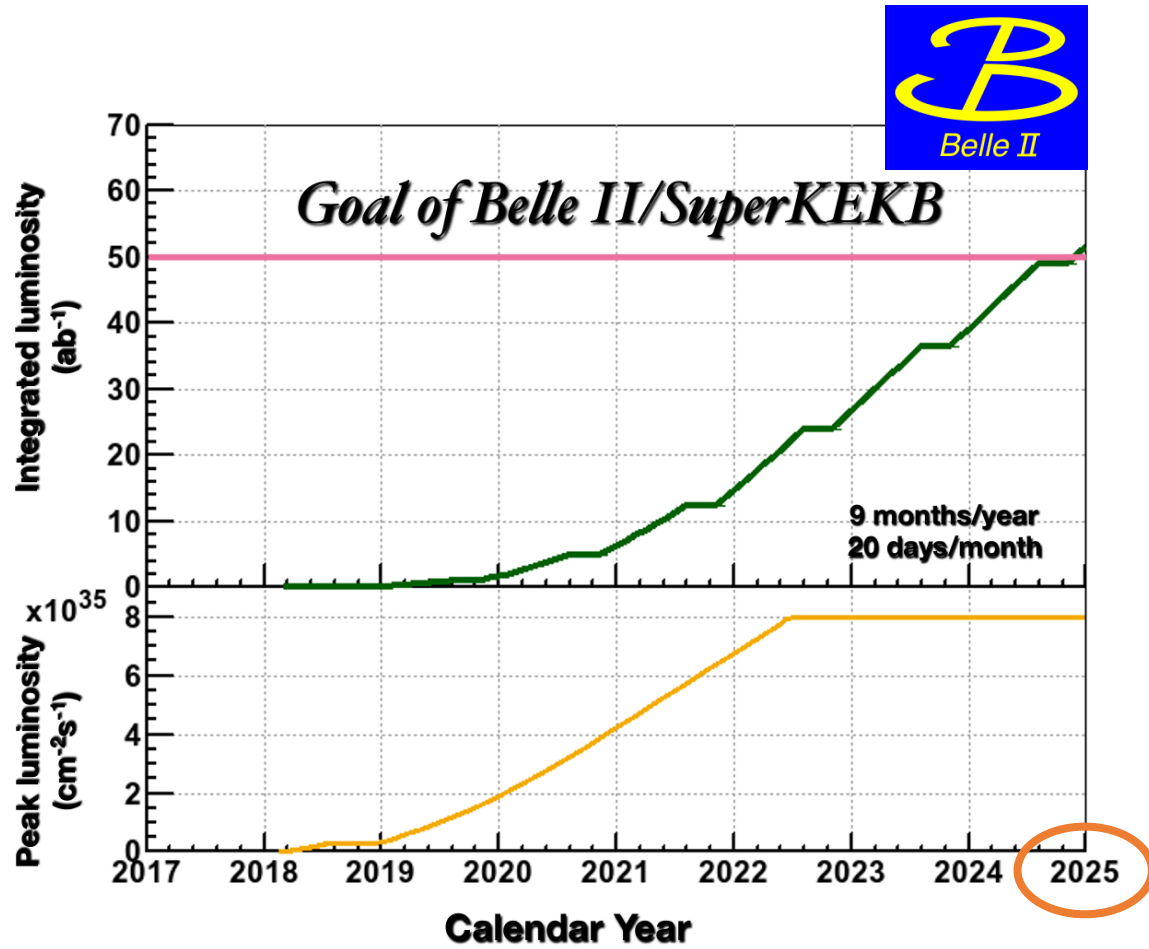
w.r.t. Belle



- HIEPA
- BINP Super  $c$ - $\tau$  factory



# Timescale



# Charm production

- $\sigma(pp \rightarrow D^0 X) @ 13 \text{ TeV} \approx 2 \text{ mb}$
- $\sigma(e^+e^- \rightarrow c\bar{c}) @ \Upsilon(4S) \approx 1.3 \text{ nb}$
- $\sigma(e^+e^- \rightarrow c\bar{c}) @ \psi(3770) \approx 6 \text{ nb}$

Parameter	Belle+BaBar (1.5 ab <sup>-1</sup> )	Belle II (50 ab <sup>-1</sup> )	LHCb (5 fb <sup>-1</sup> )	LHCb (50 fb <sup>-1</sup> )	Super $c\text{-}\tau$ (10 ab <sup>-1</sup> )
Decay time	✓		✓	✓	✗
Incoherent decays	✓			✓	✓
Coherent decays	✗			✗	✓
$N(D^0 \rightarrow K^- \pi^+) \text{ untagged}, 10^6$				40000 [1]	100
$N(D^{*+} \rightarrow D^0 \pi^+, D^0 \rightarrow K^- \pi^+), 10^6$	2.5 [2]	140 [2]	100 [1]	7000 [1]	20*
$N(D^+ \rightarrow K^- \pi^+ \pi^+), 10^6$	1.2 [3]	40	150 [1]	11000 [1]	200
$N(D_s^+ \rightarrow \varphi \pi^+), 10^6$	0.5	17	13 [1]	1000 [1]	40

\* Expected yield of  $\psi(3770) \rightarrow D^0 \bar{D}^0 \rightarrow (K^- \pi^+)(K^+ \pi^-)$  is shown for a Super  $c\text{-}\tau$  factory

[1] LHCb Collaboration, Eur. Phys. J. C73, 2373 (2013) «Implications of LHCb measurements and future prospects»

[2] Physics at Super B Factory, arXiv:1002.5012 [hep-ex]

[3] Phys. Rev. Lett. 102, 221802 (2009)

# Future precision

- An order of magnitude precision improvement can be achieved in the next decade
- The numbers shown are very approximate. Analysis of systematic uncertainties is needed

[1] “Physics at Super  $B$  Factory”, arXiv:1002.5012 [hep-ex], talk by M. Staric @ KEK FF 2014

[2] “Implications of LHCb measurements and future prospects”, Eur. Phys. J. C73, 2373 (2013)

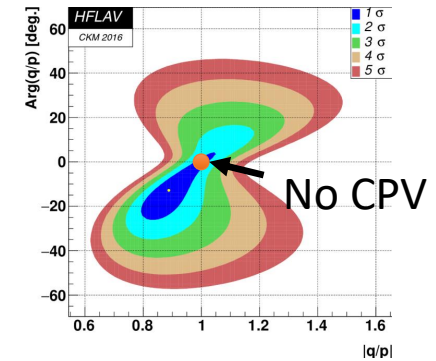
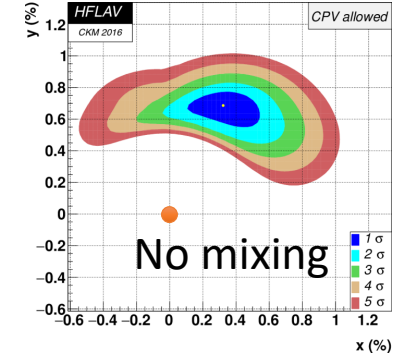
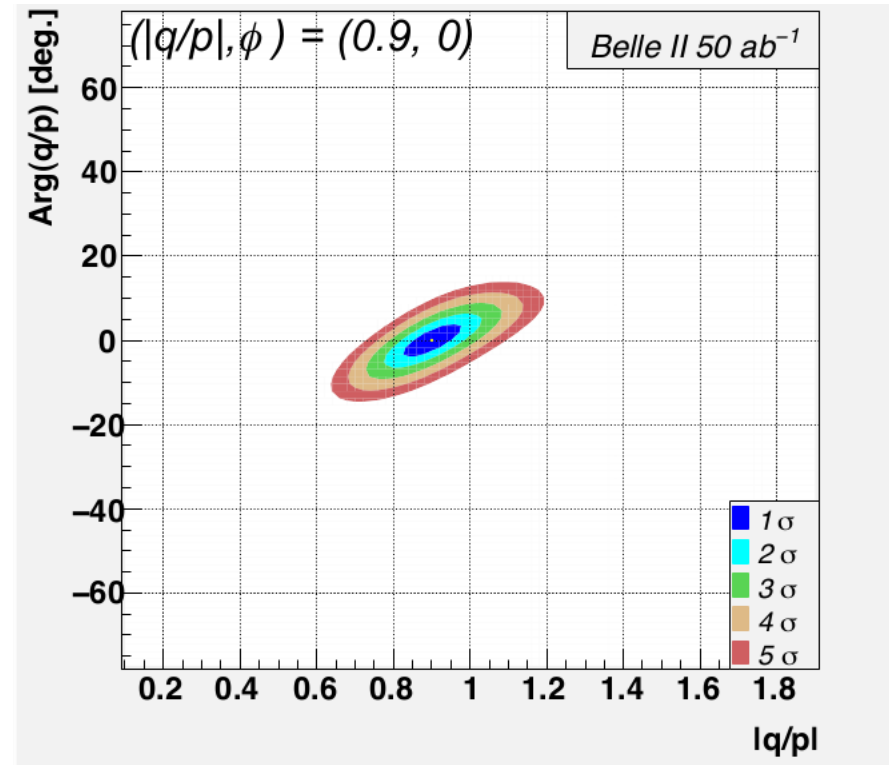
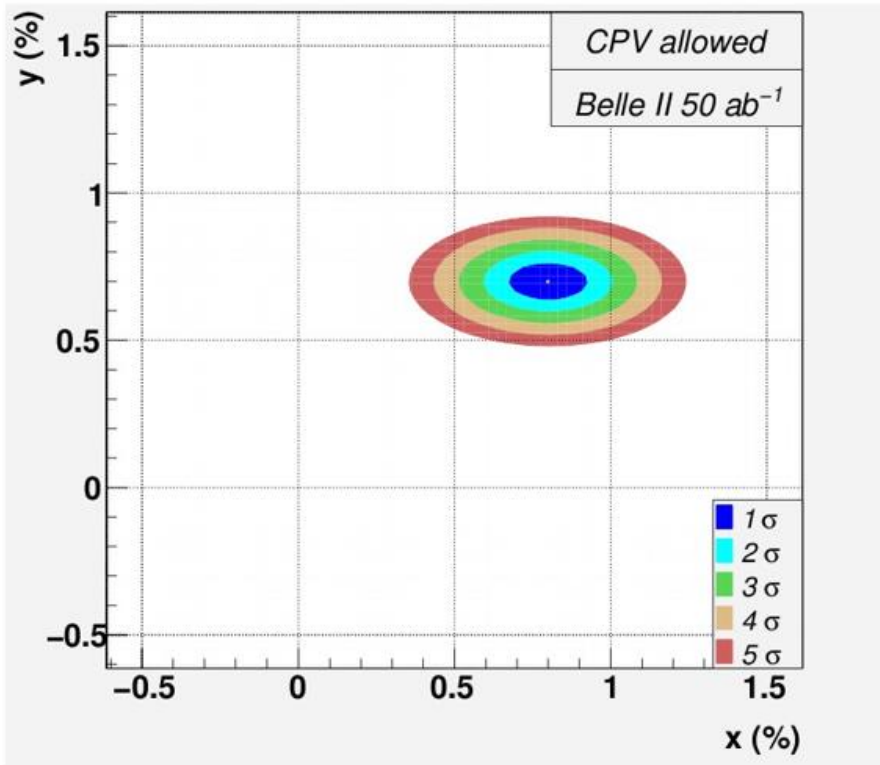
[3] A. Bondar et al., Phys. Rev. D82, 034033 (2010)

Parameter	Belle II [1] @ 50 ab <sup>-1</sup>	LHCb [2] @ 50 fb <sup>-1</sup>	Super $c\text{-}\tau$ @ 10 ab <sup>-1</sup>
WS semileptonic			
$R_M$	$\mathcal{O}(5 \times 10^{-5})$	$\mathcal{O}(5 \times 10^{-7})$	?
$D \rightarrow K\pi$ WS decays			
$y, 10^{-4}$			5
$y', 10^{-4}$	4	2	
$\cos \delta_{K\pi}$			$5 \times 10^{-3}$
$R_D, 10^{-5}$	10	0.2	1
$A_D, 10^{-4}$	3	?	?
$D \rightarrow h^+h^-$ ( $\mathcal{CP}$ eigenstates)			
$y_{\mathcal{CP}}, 10^{-4}$	4	0.4	4
$A_\Gamma, 10^{-4}$	3	$\mathcal{O}(0.1)$	?
$D^0 \rightarrow K_S^0 \pi^+ \pi^-$ Dalitz plot analysis			
$x, 10^{-4}$	8	1.7	$\mathcal{O}(1)$ [3]
$y, 10^{-4}$	5	1.9	$\mathcal{O}(1)$ [3]
$ q/p $	0.06	0.04	$\mathcal{O}(0.01)$ [3]
$\arg(q/p)$	4°	3°	$\mathcal{O}(1^\circ)$ [3]

# Global fit for Belle II @ $50 \text{ ab}^{-1}$



Alan Schwartz, IX<sup>th</sup> CKM UT Workshop (Dec. 1, 2016)

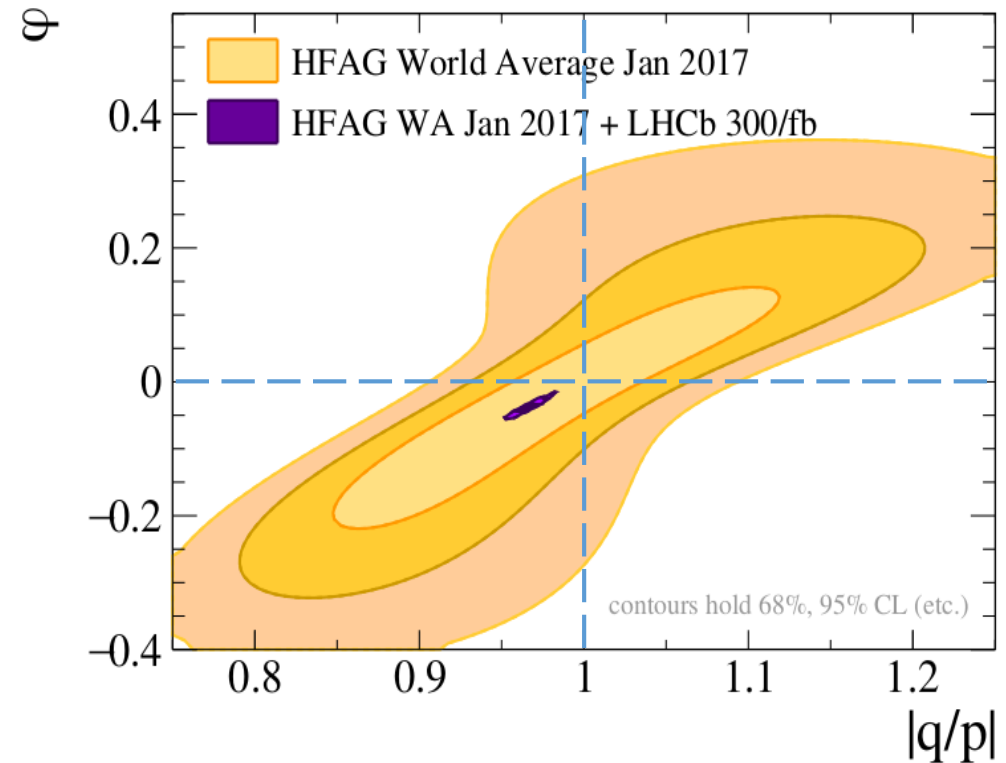
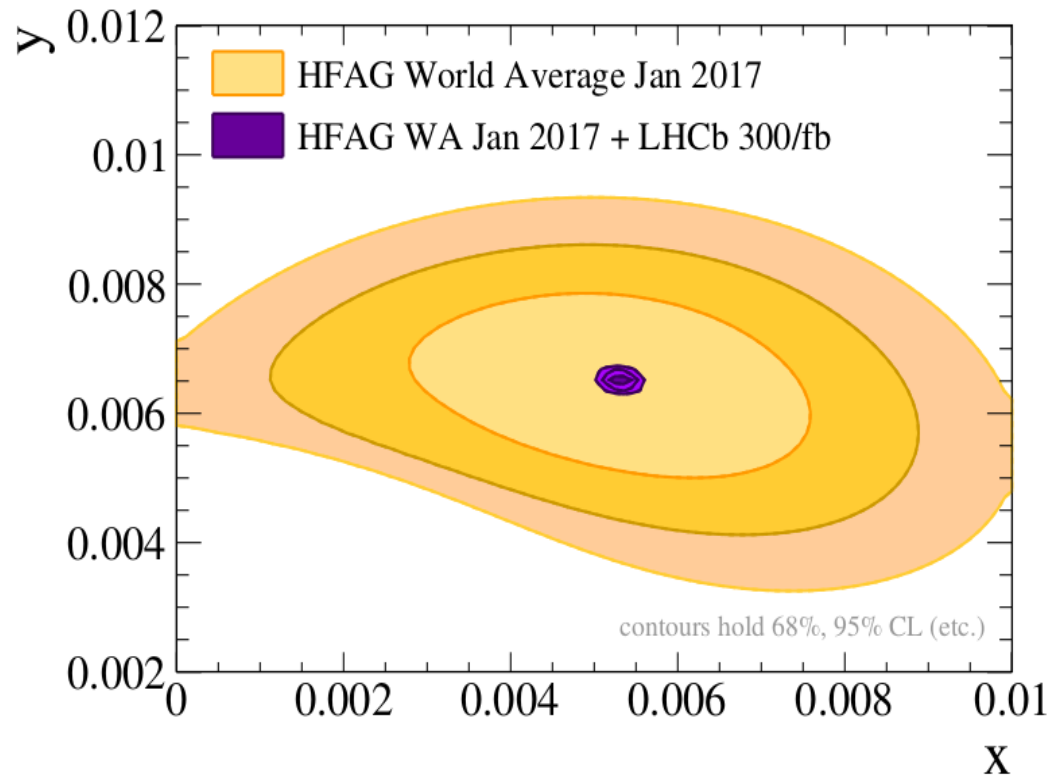


Current fit

Year 2025

# Global fit for LHCb Upgr. Phase II @ 300 fb<sup>-1</sup>

CERN-LHCC-2017-003 (Feb. 8, 2017)



Year 2035+

# Conclusions

1. Precise measurement of the charm mixing is a fundamental test of the SM
2. The existing measurements are consistent with the SM expectations
3. LHCb and Belle II are going to
  - measure the charm mixing parameters at the precision level of  $\mathcal{O}(10^{-4})$
  - access  $\mathcal{CP}$  violation in charm at the SM values ( $A_\Gamma \sim 10^{-5}$ )
4. A Super  $c$ - $\tau$  factory with  $\mathcal{L} = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  is competitive for the charm mixing measurement
5. Measurements with quantum correlations will play an important role for the future charm mixing measurements