

2 **Studies on τ decays at Belle II**

3 **L. Zani**

4 *INFN of Roma Tre,*

5 *Via della Vasca Navale, 84 Roma, Italy*

6 *E-mail:* laura.zani@roma3.infn.it

7 **ABSTRACT:** Tau leptons are powerful tools to probe physics beyond the Standard Model (SM). The
8 Belle II experiment is installed at the SuperKEKB asymmetric energy electron-positron collider
9 and aims at collecting the world's largest sample of tau pair events $e^+e^- \rightarrow \tau^+\tau^-$. Direct searches
10 for new invisible mediators, charged lepton flavor violation in τ decays, and tests of the SM via
11 precision measurements of τ lepton properties and couplings are reported in the following article.
12 The results presented here are based on the data collected by Belle II during 2019-2021.

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24 1 Introduction

25 As the only leptons massive enough to decay into hadrons, taus not only allow to investigate
26 the hadronization mechanism via their hadronic final states, but might preferentially couple to
27 non-SM physics, through mass-dependent couplings. Therefore, any possible contribution from
28 a new mediator whose coupling is proportional to lepton masses might be enhanced. From the
29 experimental point of view though, taus are challenging since they can not be detected as long-
30 lived particles, but must instead be reconstructed from their final-state products, which involve
31 undetectable neutrinos. Furthermore, they allow searching for charged lepton flavor violation
32 (LFV), which would provide an indisputable proof for beyond SM physics. Processes involving
33 LFV can occur in the SM via weak interaction charged currents, due to neutrino oscillations, and
34 are predicted at the level of 10^{-50} , which is beyond the reach of current and future experiments.
35 Belle II has a unique capability to probe both new invisible mediators and LFV in τ decays.
36 Moreover, it can look for indirect signs of non-SM physics through high precision measurements
37 of SM fundamental parameters. We report searches for new invisible particles, τ LFV decays and
38 the measurement of the τ lepton mass using the data collected by the Belle II detector [1] at the
39 SuperKEKB asymmetric energy e^+e^- collider [2]. SuperKEKB mainly operates at a centre-of-mass
40 energy (c.m.) of 10.58 GeV and adopts a nano-beam scheme to reach unprecedented instantaneous
41 luminosity. At the time of this conference, the accelerator has achieved the peak luminosity world
42 record of $4.7 \times 10^{34} \text{ cm}^{-2}/\text{s}$ and Belle II has so far collected 424 fb^{-1} of data, including unique
43 energy scan samples. It is currently in its first long shutdown.

44 **2 Leptons as discovery tools: the experimental challenges**

45 Leptonic production of tau pair processes $e^+e^- \rightarrow \tau^+\tau^-$ provide a very clean physics environment
46 and can rely on precise QED predictions to look for physics beyond the SM. The way is two-fold:
47 one could look for deviations from SM predictions in high precision measurements of very clean
48 and precisely computed observables; a second possibility is instead to search for processes that
49 would be either forbidden or highly suppressed in the SM and whose observation is *per se* a hint of
50 new physics. The first class of measurements are mainly systematically limited and to improve the
51 current results and attain the world's best precision, an excellent understanding of the experiment
52 performance at the fraction of permille level is required. On the other hand, measurements of rare
53 or forbidden processes imply to achieve unprecedented luminosity to collect sufficiently large data
54 sets and devise new analysis techniques to boost signal efficiencies in order to reach sensitivities
55 below the 10^{-8} level.

56 **3 Experimental facility: Belle II**

57 The Belle II detector, the main upgrade of its predecessor Belle, is a multipurpose spectrometer
58 surrounding the e^+e^- interaction point and providing coverage of more than 90% of the solid
59 angle. The details of the Belle II detector can be found elsewhere [1]. Belle II ensures a very
60 high reconstruction efficiency for neutral particles and excellent resolutions despite the harsh beam
61 background environment, both of which are crucial when dealing with recoiling system and missing-
62 energy final states. Additionally, it is equipped with dedicated low-multiplicity trigger lines at
63 hardware level, mainly based on calorimetric information, which were not available at Belle.
64 Profiting from the well known initial state of e^+e^- collisions, and its near-hermetic coverage, Belle
65 II has a unique capability to probe signatures involving invisible final states and particles escaping
66 detection. Moreover, the production cross-section for $e^+e^- \rightarrow \tau^+\tau^-$ events is 0.919 nb at a c.m.
67 energy $\sqrt{s} = 10.58$ GeV, allowing Belle II to collect large data samples for precision measurements
68 of τ lepton properties.

69 **3.1 Typical τ signatures in e^+e^- collisions**

70 In $e^+e^- \rightarrow \tau^+\tau^-$ processes, tau candidates are produced back-to-back in c.m. system. Their decay
71 products are well separated into two opposite hemispheres, defined by the plane perpendicular to
72 the thrust axis $\hat{\mathbf{n}}_T$, which is the vector maximizing the quantity

$$T = \max_{\hat{\mathbf{n}}_T} \left(\frac{\sum_i |\mathbf{p}_i \cdot \hat{\mathbf{n}}_T|}{\sum_i |\mathbf{p}_i|} \right), \quad (3.1)$$

73 where \mathbf{p}_i is the momentum of the final state particle i , including both charged and neutral particles.
74 According to the number of charged particles in each hemisphere, consistently with charge conser-
75 vation in τ decays, two main topologies can be selected: the 3×1 -prong decays, as schematically
76 shown in the left drawing of Figure 1, with three charged particles on one side and only one in
77 the opposite hemisphere; or the 1×1 -prong decays. Requiring the reconstructed tracks to match
78 one of these topology classes is a powerful way to suppress the main background from continuum
79 $e^+e^- \rightarrow q\bar{q}$ processes and enhance signal purity when reconstructing $e^+e^- \rightarrow \tau^+\tau^-$ events.

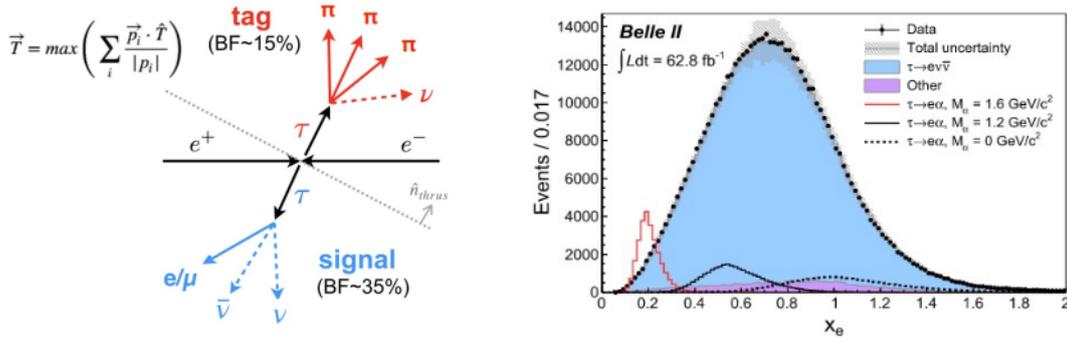


Figure 1. On the left, a typical 3x1-prong decay in $e^+e^- \rightarrow \tau^+\tau^-$ events, where one tau decays into three charged particles and the other one into one charged particle, is depicted. On the right, the distribution of the normalized lepton energy x_ℓ for the electron channel in the search for $\tau \rightarrow \ell\alpha$ is shown. Data are the black dots and the simulation is the stack filled histograms.

80 4 Searches for a new invisible boson α in τ decays

81 Decays of τ leptons to new LFV bosons are postulated in many models [3]. The process searched
 82 for in this study is $e^+e^- \rightarrow \tau^+(\rightarrow \ell\alpha)\tau^-(\rightarrow \pi^+\pi^-\pi^-\nu)$ and its charge conjugated, where the first τ
 83 is defined as the signal and the second one as the tag. The signal τ is searched for in its decay to a
 84 new invisible boson α , accompanied by a lepton $\ell = e, \mu$, therefore 3×1 -prong events are selected.
 85 The signal τ rest-frame is approximated using as energy half the collision energy $\sqrt{s}/2$ and as
 86 momentum direction the opposite to the one of the reconstructed tag τ . We exploit the kinematic
 87 features of the signal process as a two-body decay to discriminate it from the background, by looking
 88 for a narrow peak in the distribution of the normalized lepton energy in the c.m. frame, reported
 89 in the right plot of Figure 1, over a smooth contribution coming from the irreducible background
 90 of $\tau \rightarrow \ell\nu\bar{\nu}_\ell$ processes. In absence of any signal excess in 63 fb^{-1} data, 95% CL upper limits are
 91 computed on the ratio of branching fractions $\mathcal{B}(\tau \rightarrow \ell\alpha)$ normalized to $\mathcal{B}(\tau \rightarrow \ell\nu\bar{\nu}_\ell)$ [4]. This
 92 analysis provides limits between 2-14 times more stringent than the previous one set by ARGUS [5].

93 5 Direct searches for LFV $\tau \rightarrow \ell\phi$ decays

94 Possible new mediators may enhance the branching fraction for τ LFV decays $\tau^- \rightarrow \ell^-\phi$ up
 95 to observable levels of $10^{-11} - 10^{-8}$, and accommodate for flavor anomalies observed in lepton
 96 flavor universality tests with B decays [6]. In contrast to previous searches for $\tau^- \rightarrow \ell^-\phi$ decays
 97 performed at Belle [7] on $e^+e^- \rightarrow \tau^+\tau^-$ events, we apply for the first time an *untagged* approach.
 98 Only the signal τ decay to a ϕ meson candidate and a lepton, either muon or electron, is explicitly
 99 reconstructed and the other τ is not required to decay to any specific known final state. Event
 100 kinematics features and signal properties are used in a BDT classifier to suppress the background,
 101 with twice the signal efficiency for the muon mode with respect to previous analyses. Yields are
 102 extracted with a Poisson counting experiment approach from windows peaking at the known τ mass
 103 and at zero in the 2D plane of $(M_\tau, \Delta E_\tau)$, respectively, where ΔE_τ is the difference between the
 104 reconstructed energy of the signal τ in the c.m. frame and half the collision energy. We find no

105 significant excess and set 90% CL upper limits on the branching fractions of $\mathcal{B}_{\text{UL}}(\tau \rightarrow e\phi) =$
 106 23×10^{-8} and $\mathcal{B}_{\text{UL}}(\tau \rightarrow \mu\phi) = 9.7 \times 10^{-8}$ [8].

107 6 Measurement of the τ lepton mass

108 Lepton properties are fundamental parameters of the SM and need to be measured with the highest
 109 precision. Belle II is suitable to measure several τ lepton properties. By applying the pseudo-
 110 mass M_{min} technique to reconstructed $e^+e^- \rightarrow \tau^+\tau^-$ events from 190 fb^{-1} data, we provide the
 111 world's most precise measurement of the τ mass M_τ . The measured value is extracted from a fit to
 112 the endpoint of the distribution $M_{\text{min}} = \sqrt{M_{3\pi}^2 + 2(\sqrt{s}/2 - E_{3\pi}^*)(E_{3\pi}^* - P_{3\pi}^*)}$, which is computed
 113 from events where the signal τ is reconstructed in its decays to three charged pions and the other τ
 114 decaying into one charged particle. The distributions of the pseudomass in simulation and data is
 115 shown in the left plot of Figure 2. An excellent control of the systematic sources, dominated by the
 116 calibration of the beam energies and the charged-particle momentum scale, is required to reduce
 117 the total systematic uncertainty to $0.11 \text{ MeV}/c^2$, achieving the most precise measurement to date of
 the τ lepton mass of $1777.09 \pm 0.08_{\text{stat}} \pm 0.11_{\text{sys}}$ [9].

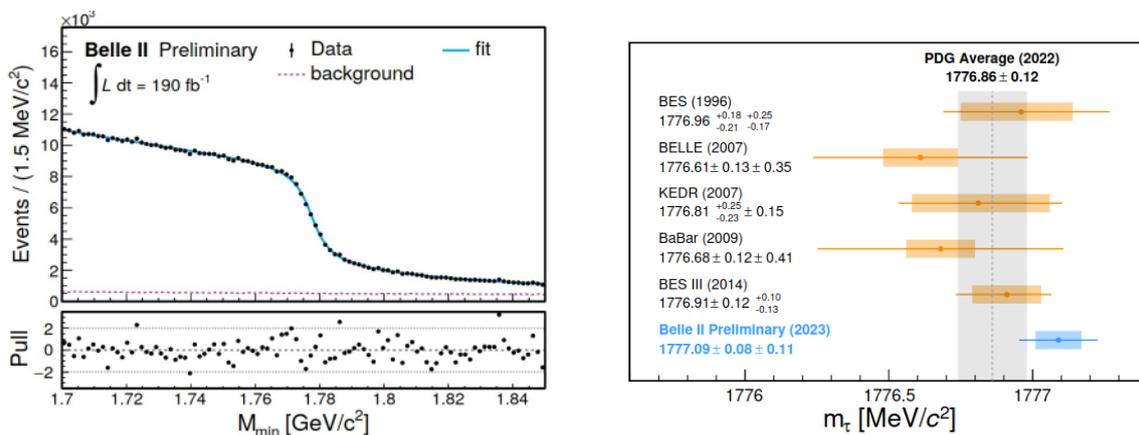


Figure 2. On the left, the spectrum of the reconstructed pseudomass in data (black dots) and the superimposed fit (solid blue line) are shown. The bottom inset plot displays differences between data and fit result divided by the statistical uncertainties. On the right, the summary of the most precise measurements of the tau mass to date, compared to the world average (gray band) and this work result (blue text).

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119 7 Prospects on lepton flavor universality tests

120 Lepton flavor universality (LFU) in SM assumes all three leptons have equal coupling strength to
 121 the charged gauge bosons of the electroweak interaction. Many models predict new forces violating
 122 LFU, that could for example imply a new singly-charged scalar singlet [10]. Tau decays allow high
 123 precision tests of LFU by measuring the ratio of the branching fractions of τ decays to muon and to
 124 electron,

$$R_\mu = \frac{\mathcal{B}(\tau \rightarrow \mu\nu_\mu\nu_\tau)}{\mathcal{B}(\tau \rightarrow e\nu_e\nu_\tau)} \quad (7.1)$$

125 The most precise result to date is $R_\mu = 0.9796 \pm 0.0016_{\text{stat}} \pm 0.0036_{\text{sys}}$ provided by Babar [11]. It uses
 126 467 fb^{-1} collision data, for a final 0.4% precision, systematically dominated by the contribution of
 127 the particle identification (PID) and trigger uncertainties. Simulation studies at Belle II show room
 128 for a several improvements: by devising dedicated low multiplicity triggers based on calorimeter
 129 information, which provide a better understanding of the kinematic dependency and reduce the
 130 associated systematic uncertainty; by dropping the likelihood-based PID selector for pions and
 131 deploying BDT classifier for lepton identification, which will decrease the probability to wrongly
 132 identify a pion as a lepton to less than 0.1%; eventually, adding the 1×1 -prong decays as signal
 133 signature to increase the size of the analyzed data set. Further studies for the development of the
 134 specific 1×1 topology triggers are still needed, but already with one quarter of Babar data set, Belle
 135 II expected sensitivity achieves the same statistical precision of 0.16%.

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